

Nanowire-based superconducting qubits

L. Y. Cheung, R. Haller, J. H. Ungerer, C. Ciaccia, A. Kononov, H. Zheng, N. Sangwan, T. Jenniskens T. Kanne, J. Nygard, P. Winkel, T. Reisinger, I. M. Pop, E. P.A.M. Bakkers, J. Ridderbos, A. Baumgartner, Christian Schönenberger

Workshop on Superconductor-Semiconductor Hybrids
March 12-14 2024
Niels Bohr Institute / The University of Copenhagen
by Christian Schönenberger

Quantum- and Nanoelectronics group: www.nanoelectronics.ch
Swiss Nanoscience Institute: <https://nanoscience.ch/en/>
Physics Department of the University of Basel: <https://physik.unibas.ch>

Nanowire-based superconducting qubits

Han Zheng



Luk Yi Cheung



Nikunj Sangwan
EPFL



Tom Jenniskens
TUE



Artem Kononov



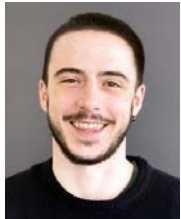
Roy Haller



Joost Riderbos



Carlo Ciaccia



Jann Ungerer



Andreas Baumgartner



Christian Schönenberger



acknowledgement:

Group **Jesper Nygard** for continuous support and collaboration



Niels Bohr Institute

Group **Erik Bakkers**, growth of *Ge/Si nanowires*



Univ. of Lund team: **K. Dick**,
C. Thelander and **V. Maisi**



Joost Ridderbos (present address Univ. of Twente)

UNIVERSITY
OF TWENTE.

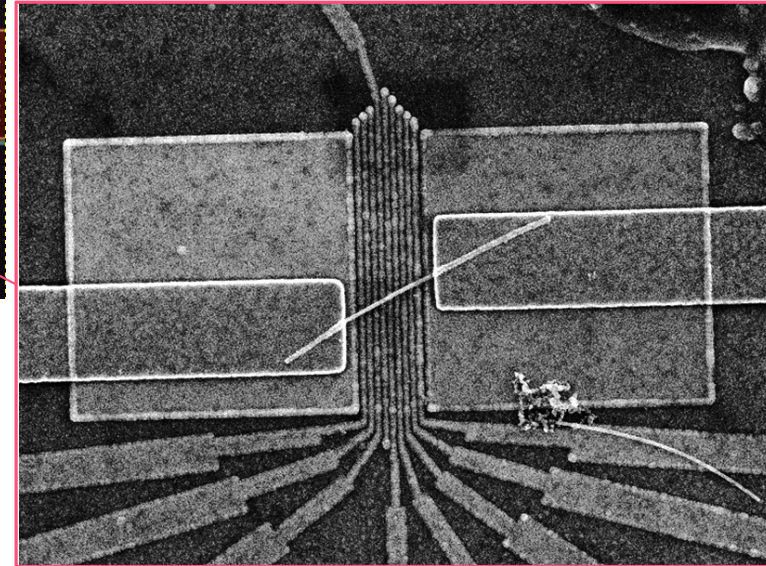
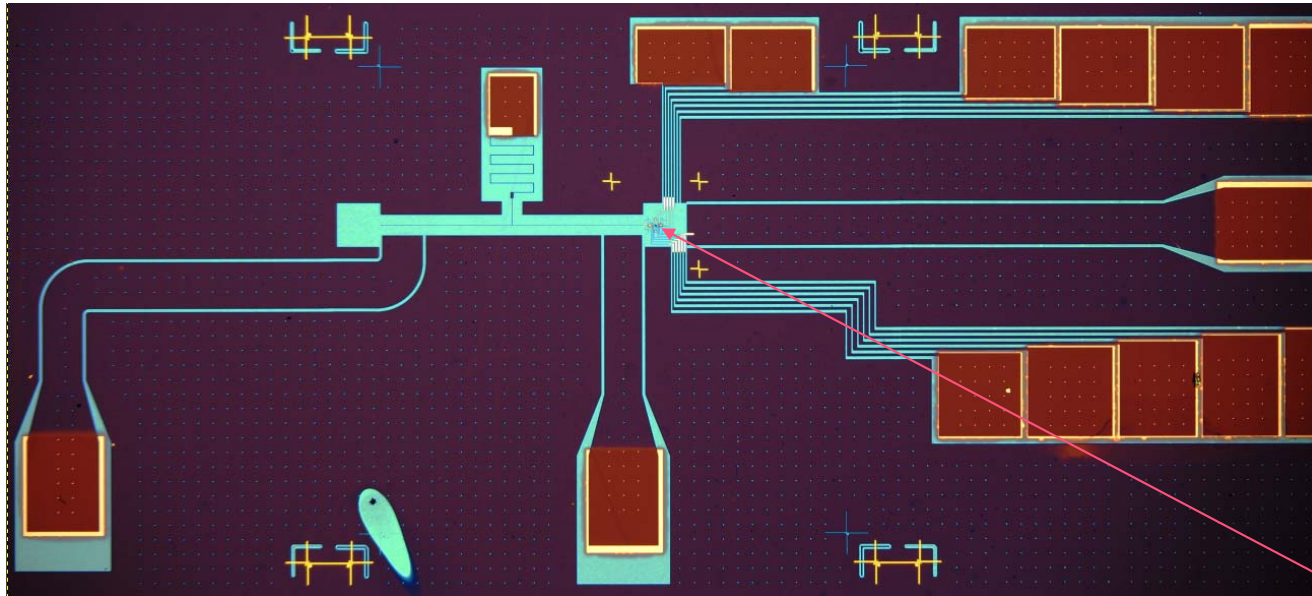
Group **Ioan Pop**, *parametric amplifier*



Nanowire-based qubit research at CS's group @ University of Basel

Nanowire-based qubit research at CS's group @ University of Basel

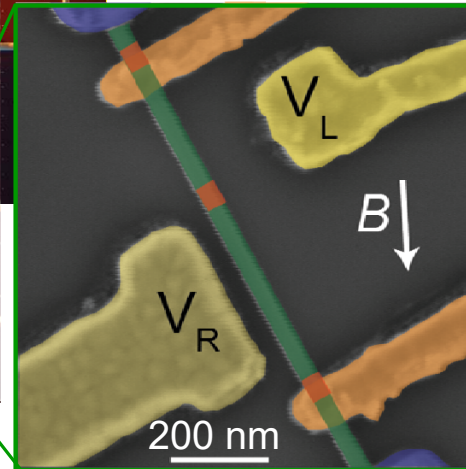
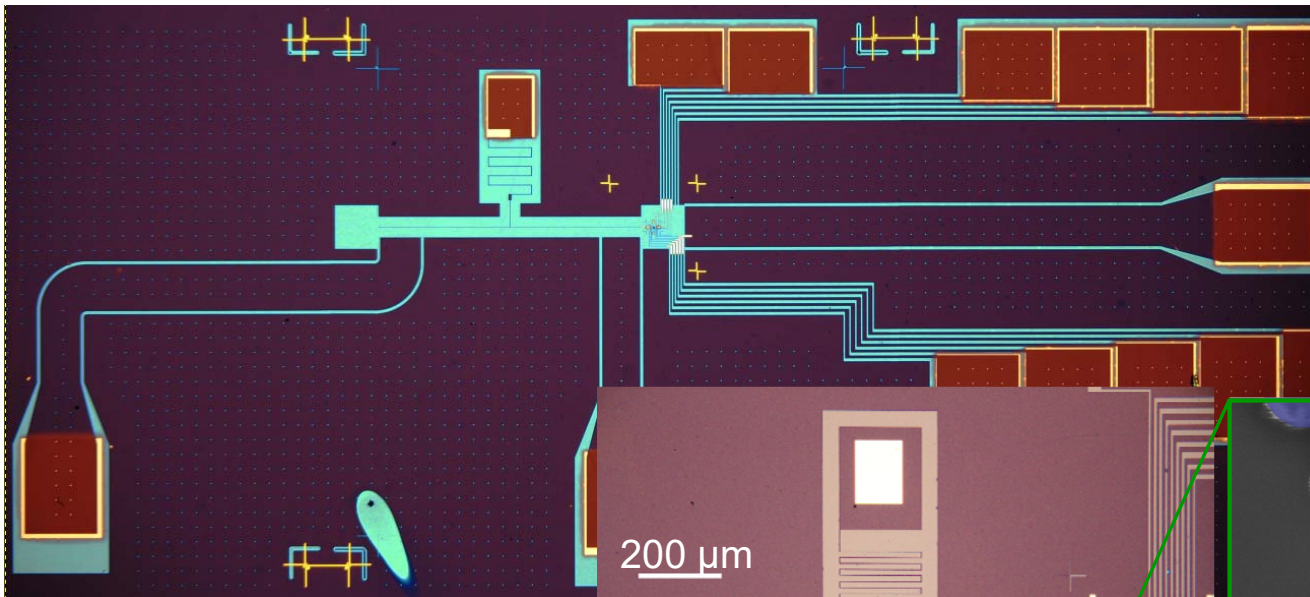
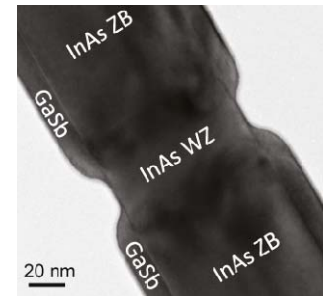
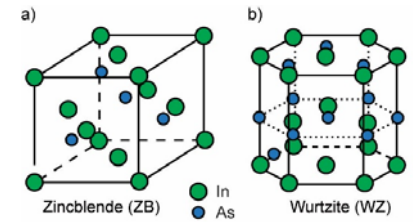
Spin qubits in GeSi core shell nanowires (based on double quantum dots) with Dominik Zumbühl's group



NbTiN high-impedance $\lambda/2$ cavity
readout in transmission

Nanowire-based qubit research at CS's group @ University of Basel

Spin qubits in InAs NWs (based on double quantum dots)



excellent NWs from Lund yielding very stable double quantum dot devices

NbTiN high-impedance $\lambda/2$ cavity readout in transmission

Strong coupling between a microwave photon and a **singlet-triplet qubit**,

J. H. Ungerer et al. Nature Comm. 15, 1068 (2024)

Nanowire-based qubit research at CS's group @ University of Basel

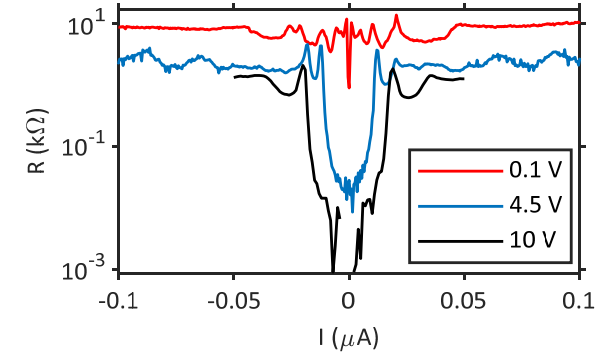
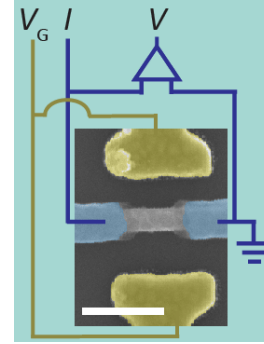
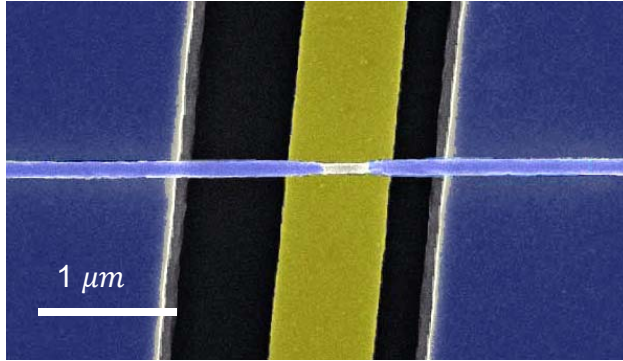
- **Spin qubits** in GeSi core shell nanowires (based on double quantum dots)
- **Spin qubits** in InAs NWs (based on double quantum dots)

Nanowire-based qubit research at CS's group @ University of Basel

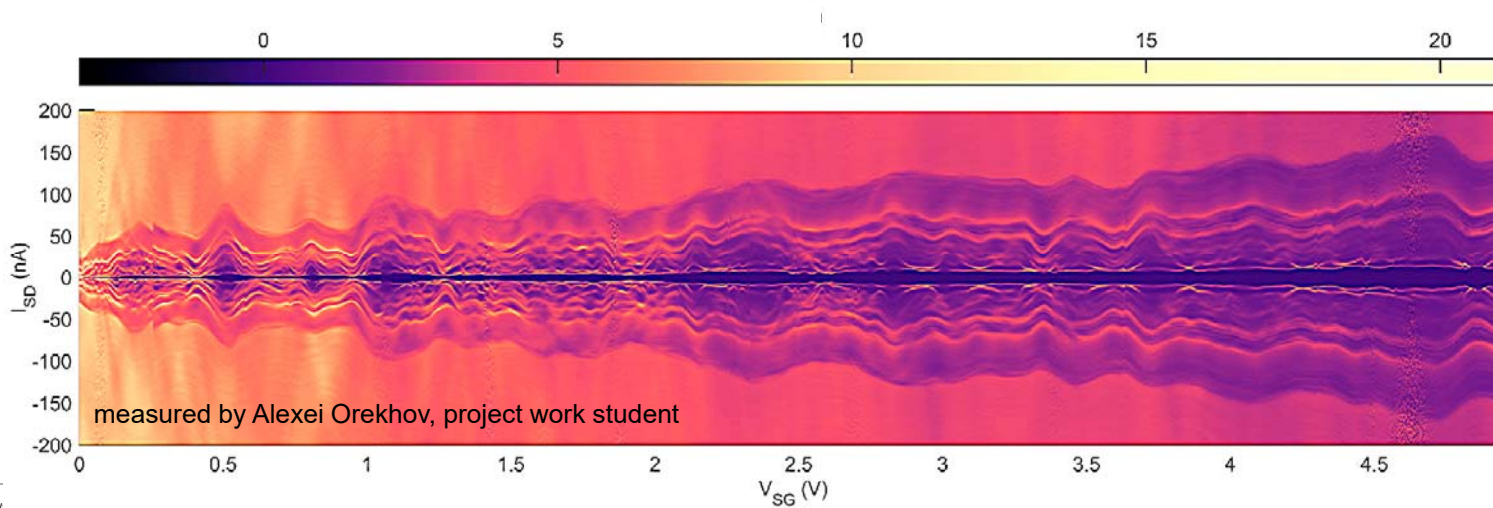
- **Spin qubits** in GeSi core shell nanowires (based on double quantum dots)
- **Spin qubits** in InAs NWs (based on double quantum dots)
- **Andreev (spin) qubits** in InAs nanowires (material from Jesper Nygard's group)
- **Gatemon qubits** in GeSi core shell nanowires (material from Erik Bakkers' group)

InAs nanowires with in-situ grown Al shell (few facets to full shell)

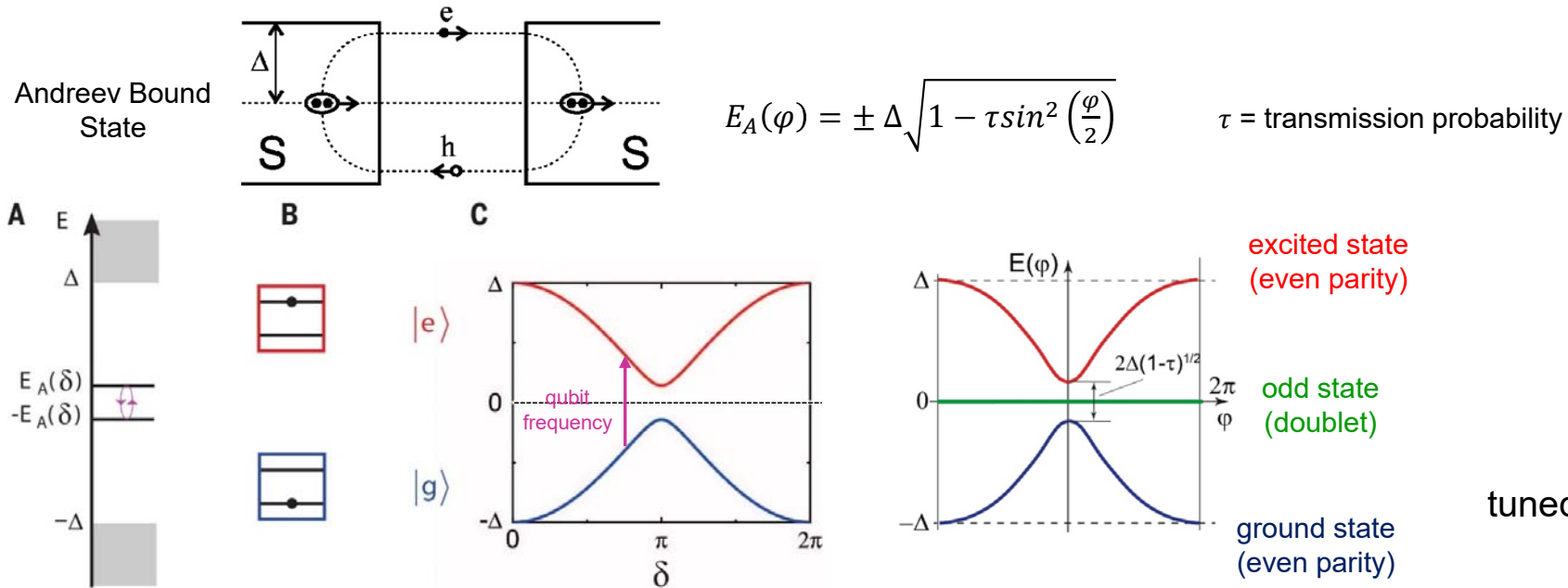
InAs full (Al) shell nanowires, obtained from **Jesper Nygard** (CPH)
Etched and **gated** device. Here, the NW is suspended over the gate



qubit
tuned by gate voltage



The Andreev Qubit



qubit
tuned by phase bias

see for example: M. F. Goffman et al.
Supercurrent in atomic point contacts
and Andreev states, Phys. Rev. Lett. 85,
170 (2000).

L. Bretheau et al. Exciting Andreev pairs
in a superconducting atomic contact,
Nature 499 (7458), 312-315 (2013).

The AndQC FETopen project

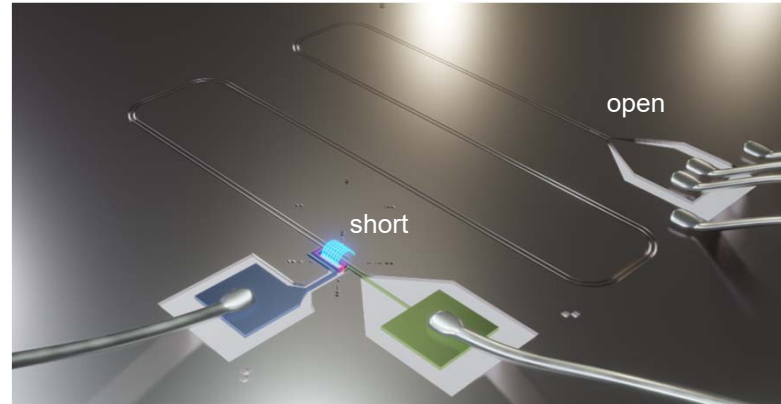


Andreev qubits for scalable quantum computation

consortium partner

- Chalmers
- TU Delft
- BME
- CPH
- CNR
- UAM
- CEA
- UBAS

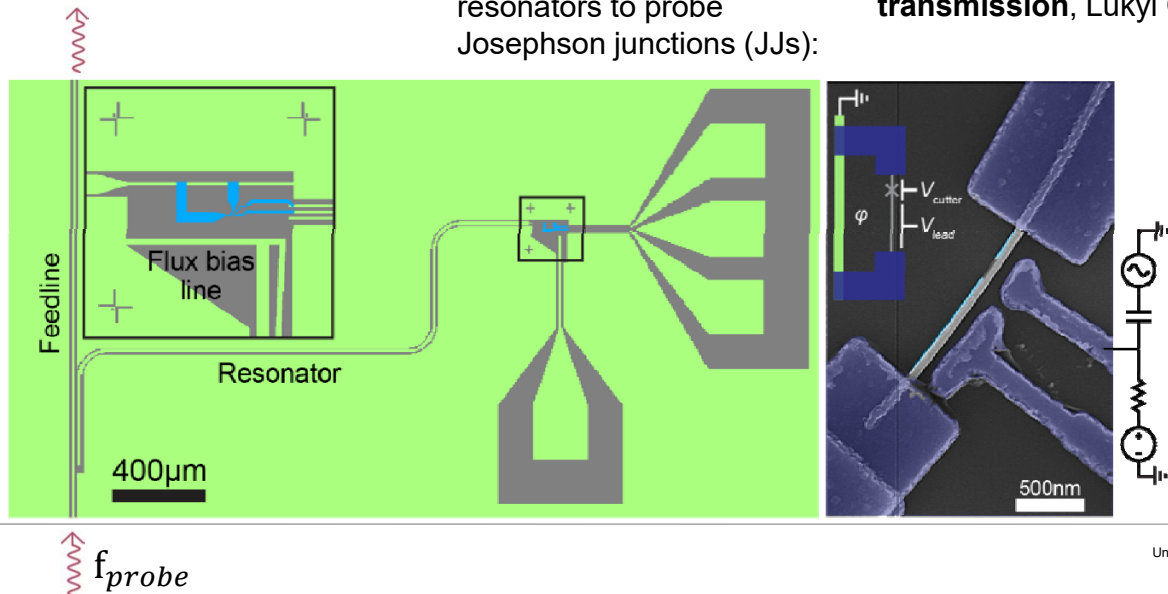
Our goal is to establish the foundations of a **radically new** solid state platform for scalable quantum computation, based on **Andreev qubits**. To carry out this research program, we rely on the instrumental combination of experimentalists, theorists and material growers, together having the necessary expertise on all aspects of the proposed research



1st generation of rf resonators to probe Josephson junctions (JJs): $\lambda/4$ resonator in **reflection**, illustration by Roy Haller

2nd generation of rf resonators to probe Josephson junctions (JJs):

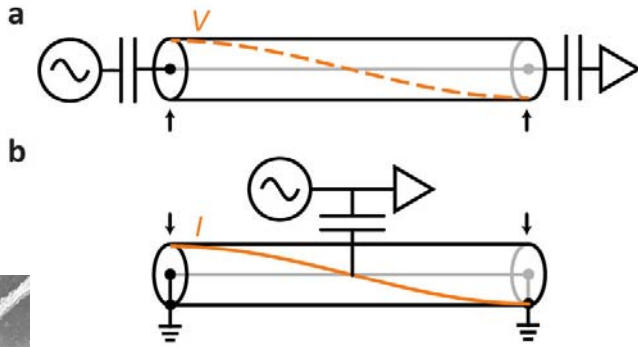
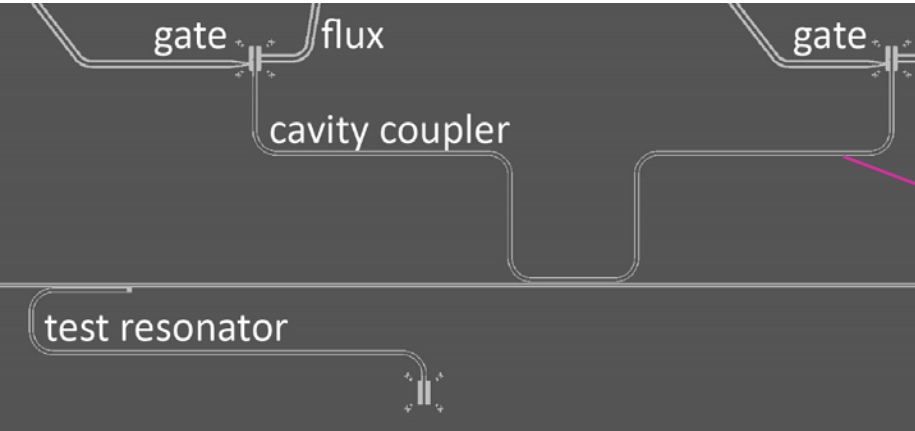
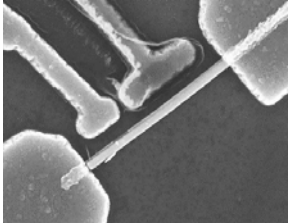
$\lambda/4$ resonator with **feedline in transmission**, Lukyi Cheung et al



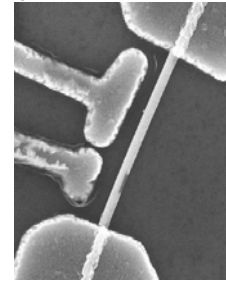
Remote Qubit-Qubit coupling



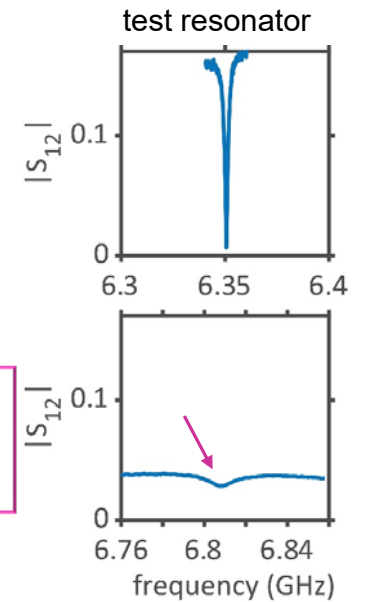
qubit L



qubit R

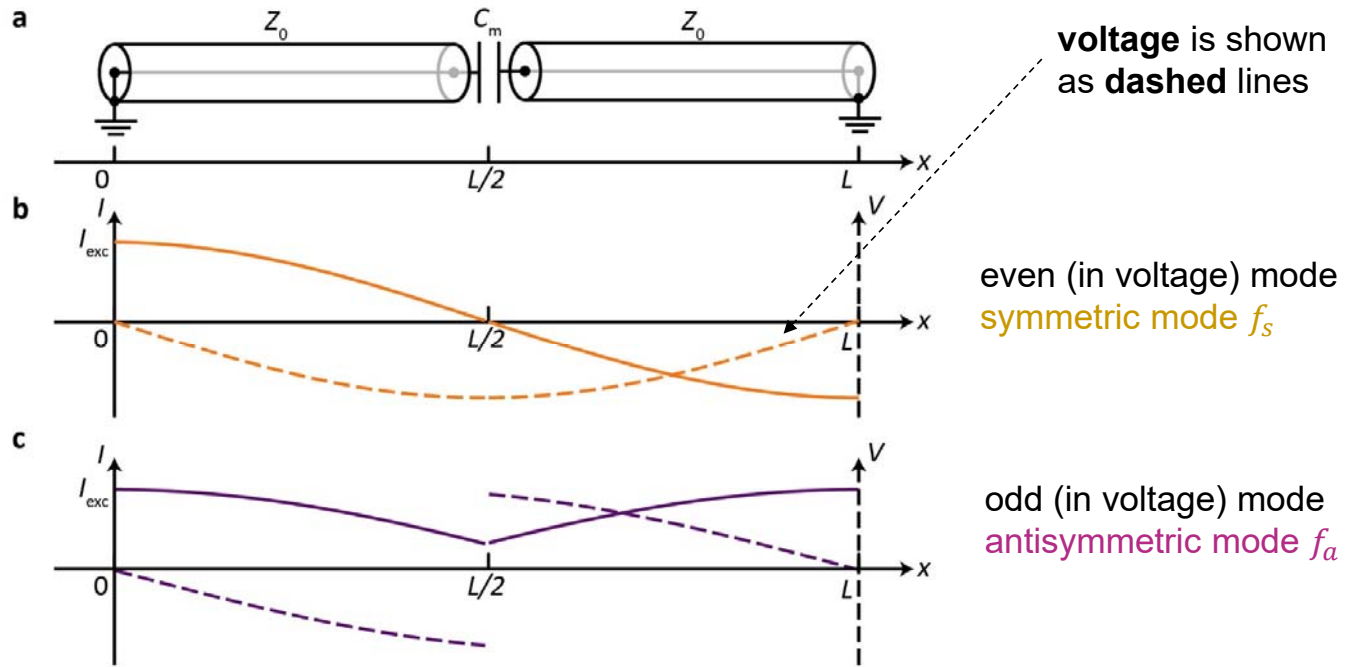


cavity coupler shallow dip ?
 $Q_i < 10^3$



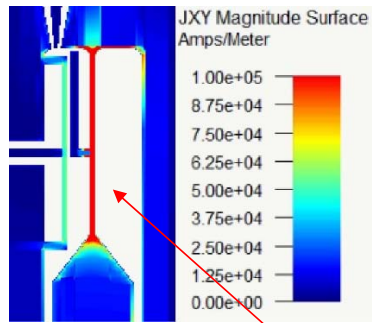
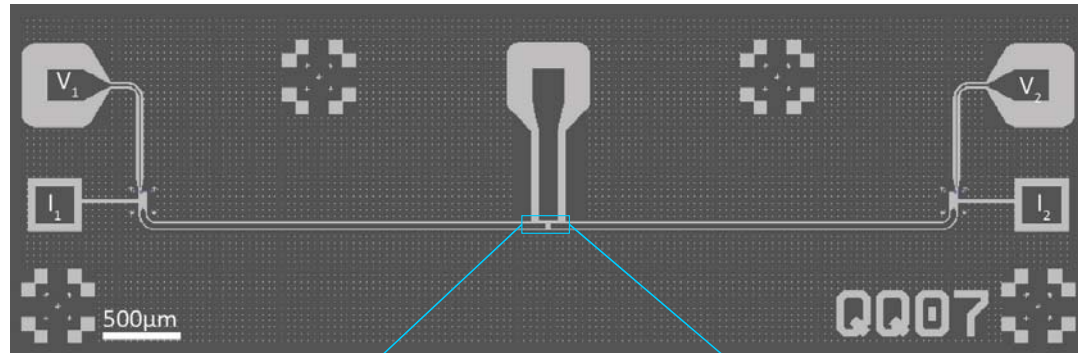
Remote Qubit-Qubit coupling

cut the "ground loop"
2 coupled $\lambda/4$ resonators

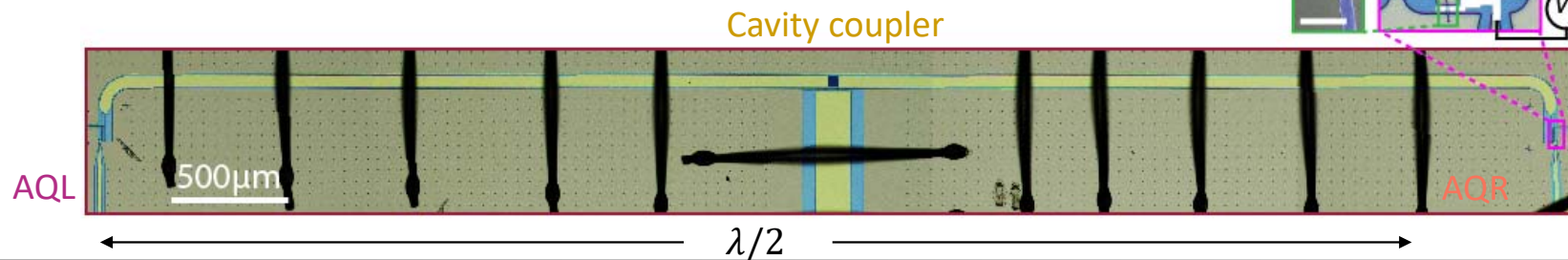
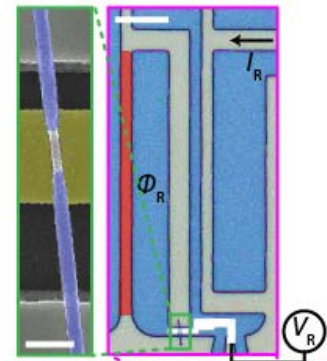


design rule: $f_s - f_a \sim 5g \cong 500$ MHz

Remote Qubit-Qubit coupling

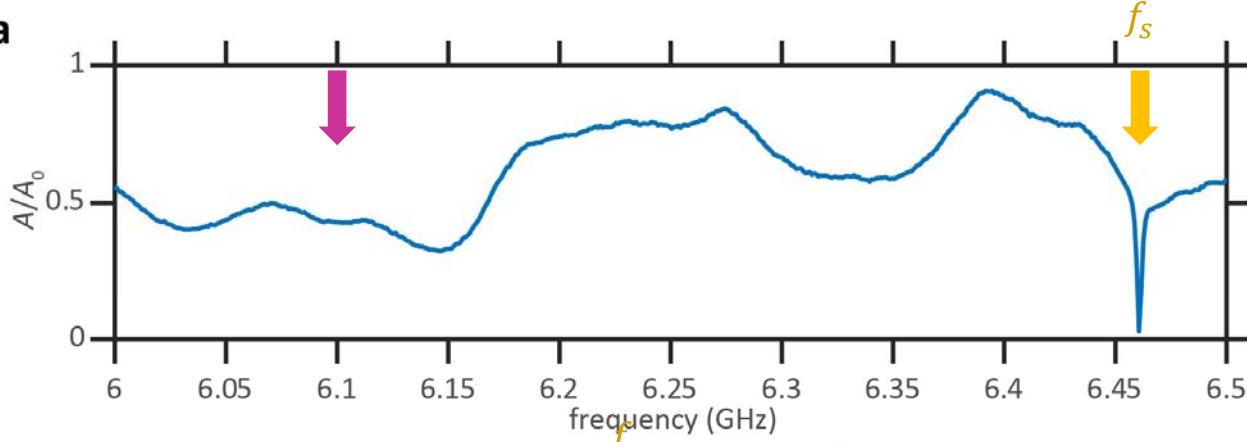


shared inductor $l \cong 90 \text{ pH}$

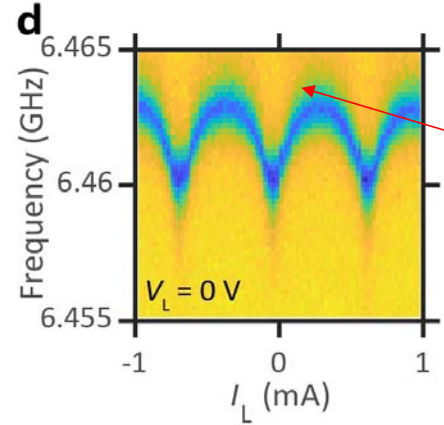
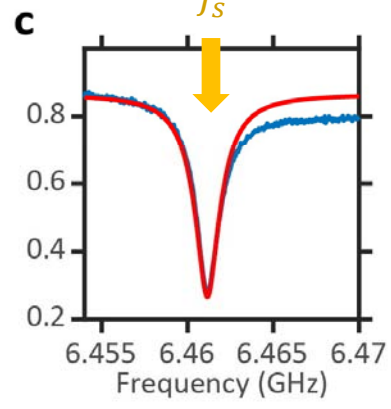
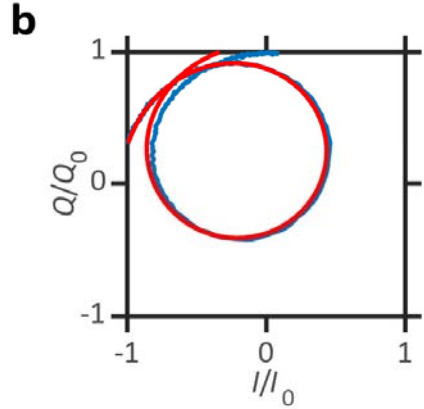


Remote Qubit-Qubit coupling

Resonator characterization:

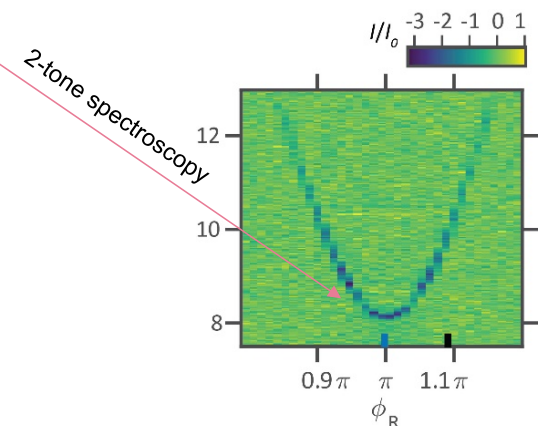
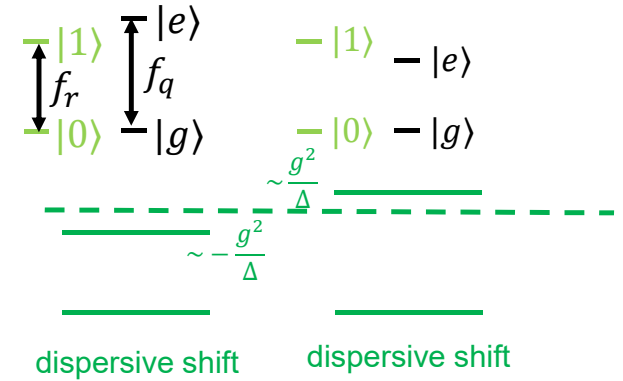
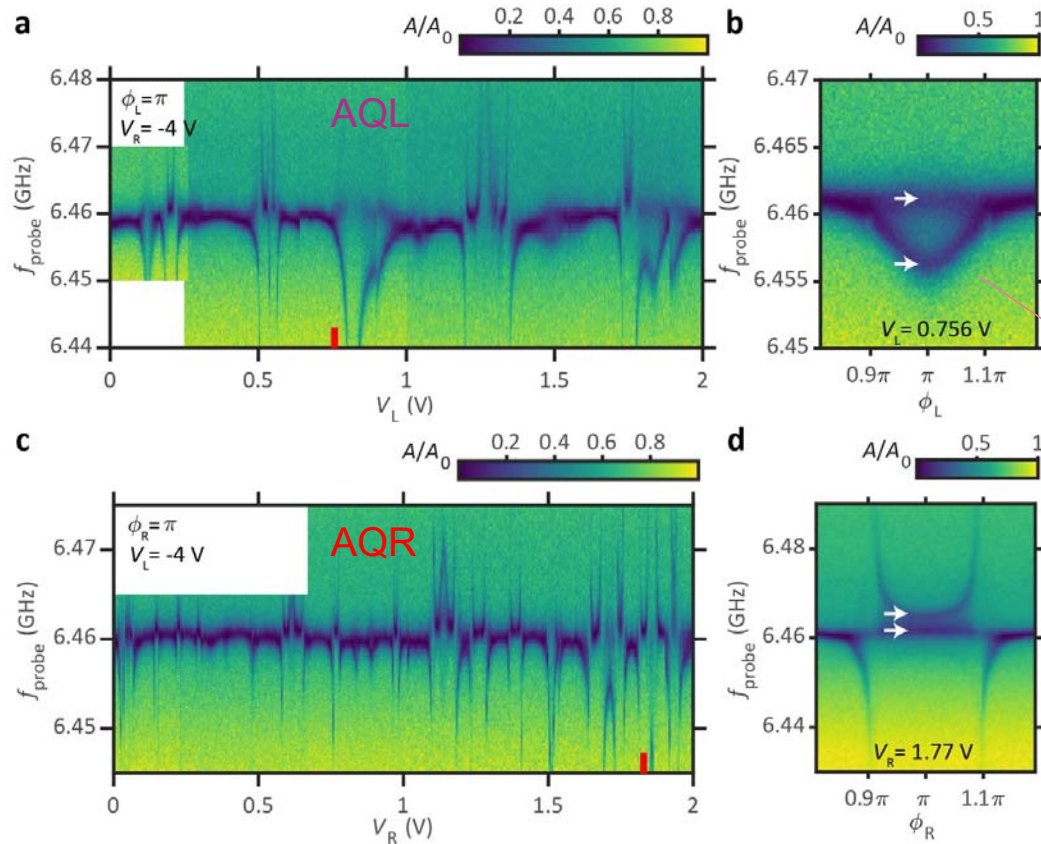


$\kappa_i/2\pi = 0.7$ MHz
 $\kappa_{CS}/2\pi = 1.2$ MHz



periodic dispersive shift due to AQL being phase-tuned by the flux-bias given in current

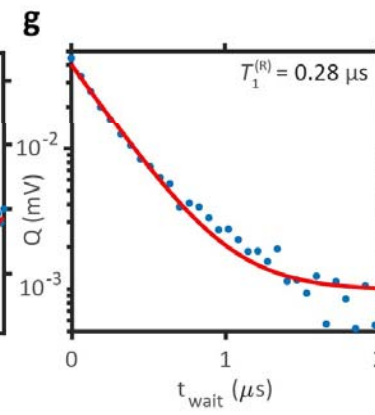
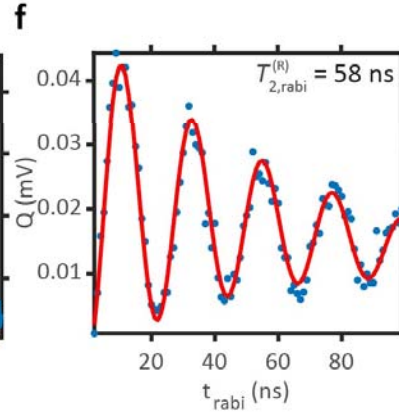
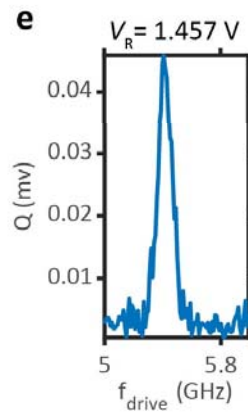
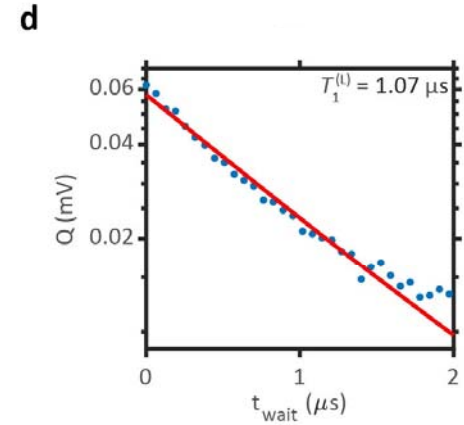
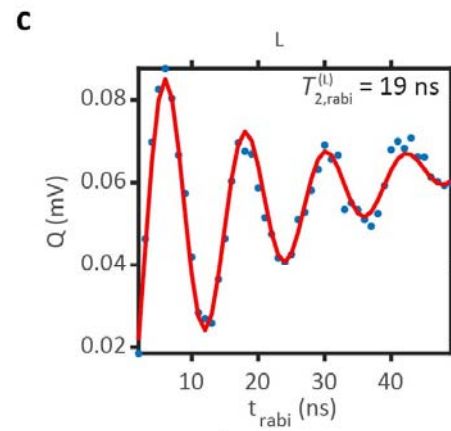
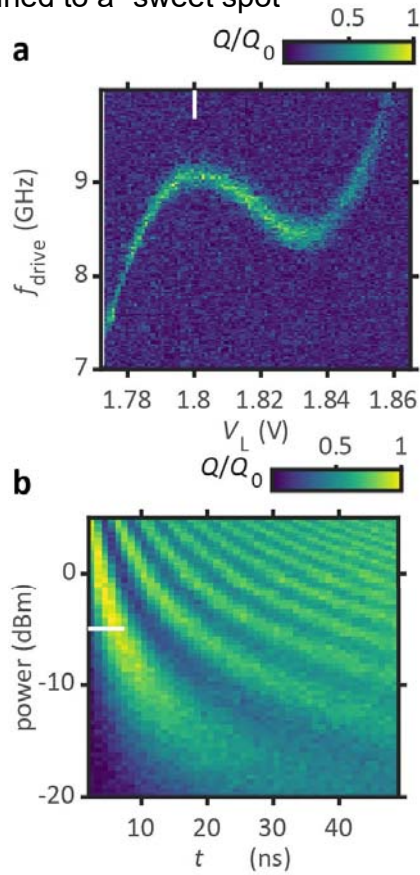
Interaction of AQL & AQR with f_s



dispersive shift $\frac{g^2}{\Delta} \cong 5.5 \text{ MHz}$
 with a detuning $\Delta \cong 1.8 \text{ GHz}$, yields
 a coupling strength $g = 100 \text{ MHz}$

Coherence of Single Qubits

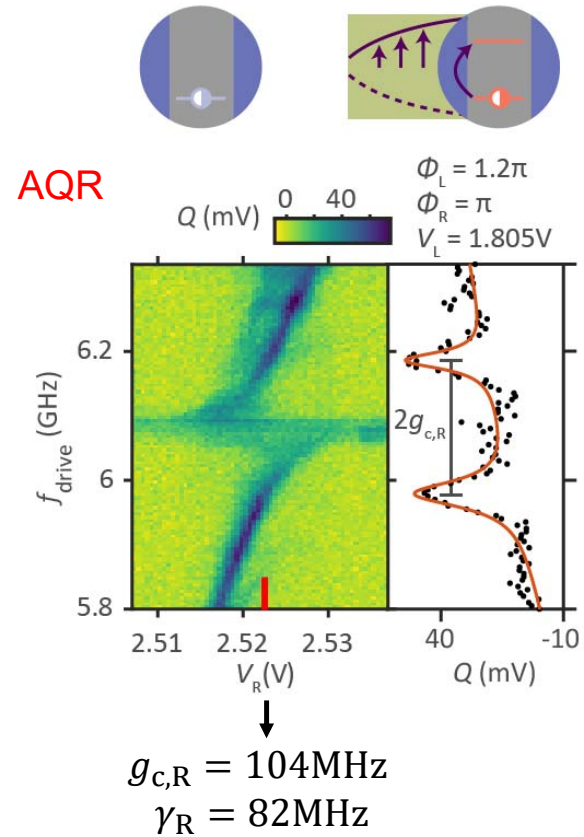
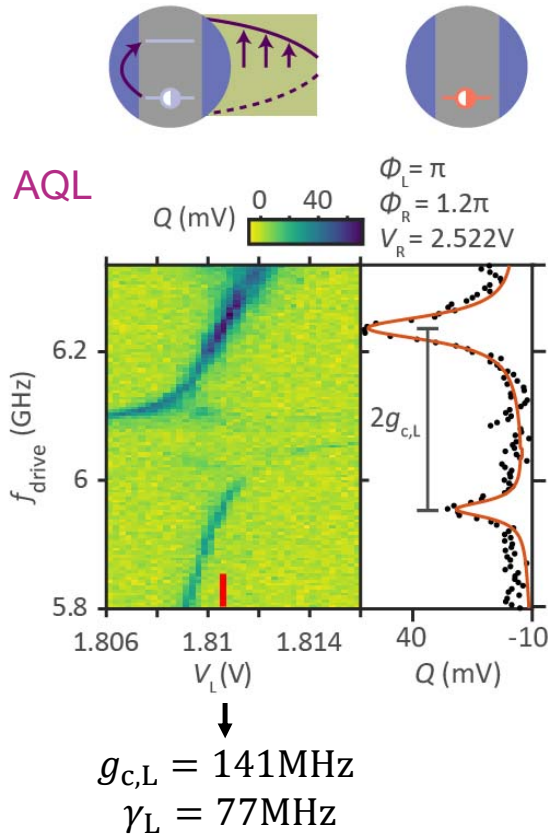
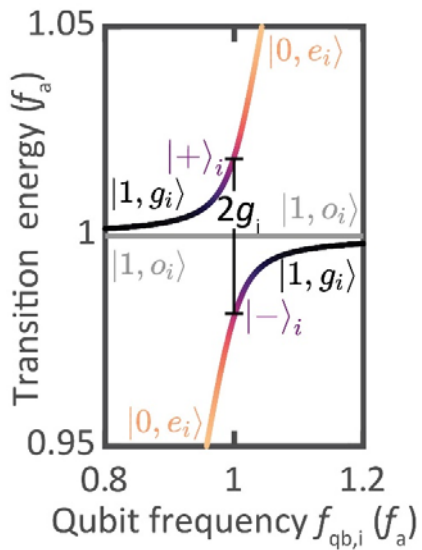
AQL tuned to a "sweet spot"



$$T_1 \cong 1 \mu\text{s} \text{ and } T_{2,\text{rabi}} \sim 50 \text{ ns}$$

Qubit Coupling to the Asymmetric Mode f_a

Pulsed differential two-tone spectroscopy at $f \sim f_a$ with a probe tone at $f \sim f_s$



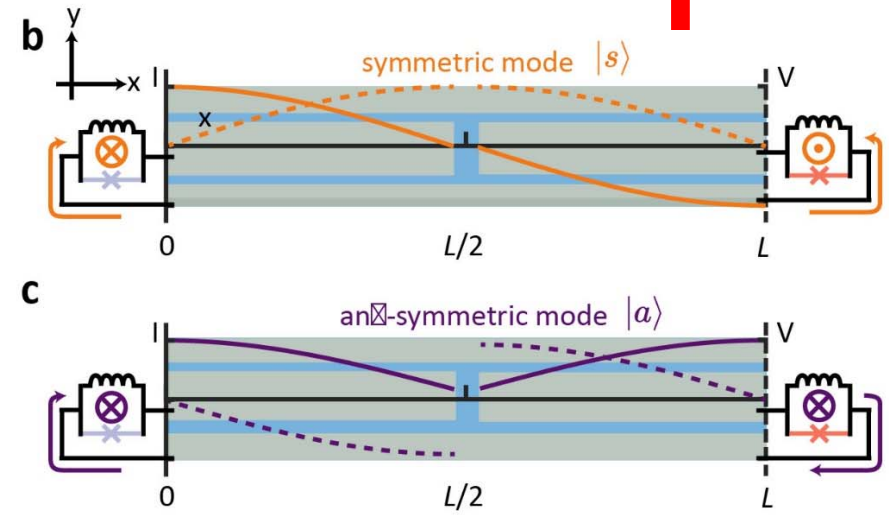
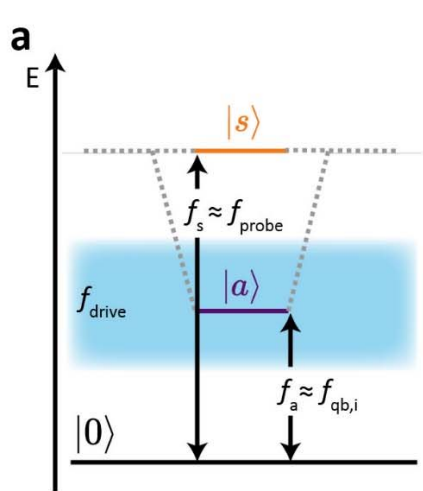
Sign of Qubit-Qubit Coupling depends on Symmetry

Tavis-Cummings Hamiltonian

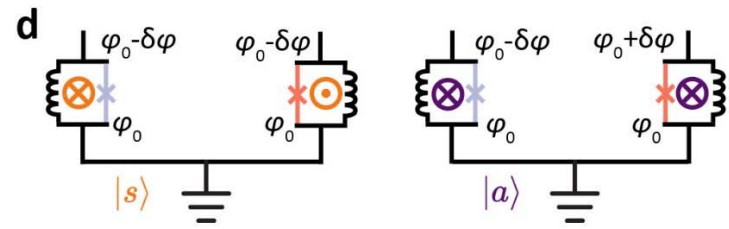
$$H_{TC} = \frac{\hbar f_{qb,L}}{2} \sigma_{z,L} + \hbar f_a (a^\dagger a) + \frac{\hbar f_{qb,R}}{2} \sigma_{z,R} + \hbar g_{c,L} (a \sigma_{+,L} + a^\dagger \sigma_{-,L}) - \hbar g_{c,R} (a \sigma_{+,R} + a^\dagger \sigma_{-,R})$$



conceptional idea:
use asymmetric mode
as qubit-qubit coupler
mode and the symmetric
one for qubit readout

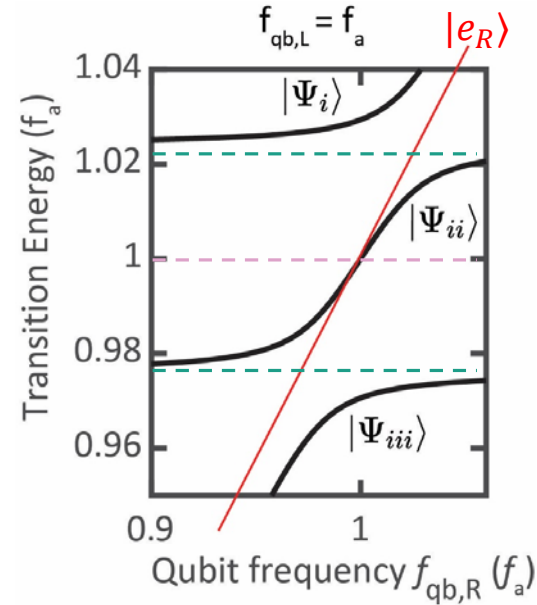
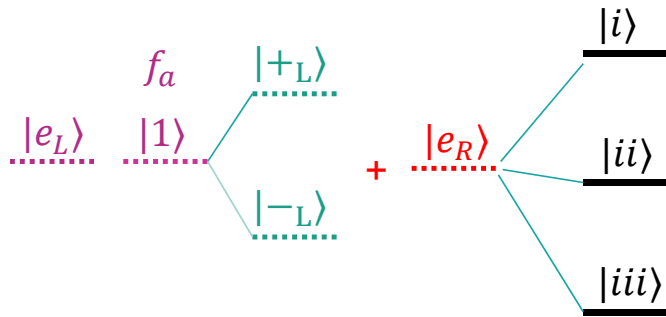


symmetric phase response



antisymmetric phase response

Strong Qubit-Qubit-Resonator Coupling



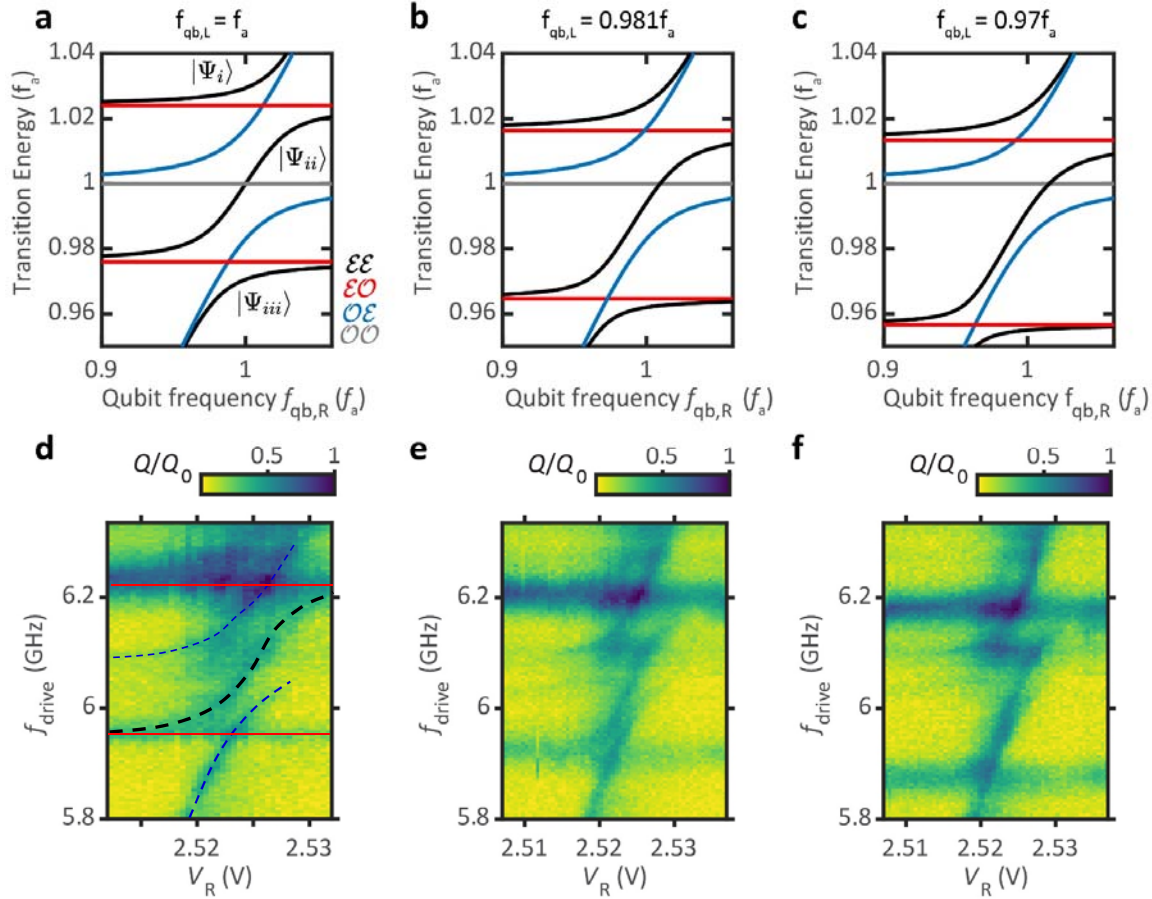
dressed states under resonance condition:

$$|\Psi_i\rangle \rightarrow |\Psi_s\rangle = \frac{1}{\sqrt{2}} (|1, g_L, g_R\rangle + \frac{1}{\sqrt{2}} (|0, e_L, g_R\rangle + |0, g_L, e_R\rangle)) \quad (7.7)$$

$$|\Psi_{ii}\rangle \rightarrow |\Psi_q\rangle = \frac{1}{\sqrt{2}} (|0, e_L, g_R\rangle - |0, g_L, e_R\rangle) \quad (7.8)$$

$$|\Psi_{iii}\rangle \rightarrow |\Psi_a\rangle = \frac{1}{\sqrt{2}} (|1, g_L, g_R\rangle - \frac{1}{\sqrt{2}} (|0, e_L, g_R\rangle + |0, g_L, e_R\rangle)). \quad (7.9)$$

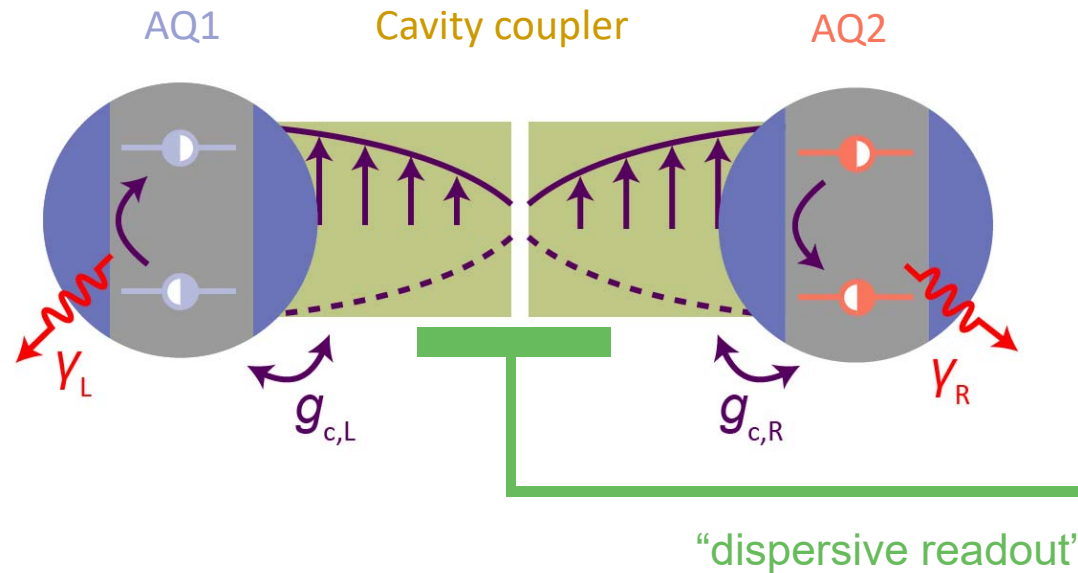
Strong Qubit-Qubit-Resonator Coupling



$|\psi_i\rangle$ & $|\psi_{iii}\rangle$ barely visible, why? \rightarrow Symmetry of coupling

(control experiment with symmetric mode)

Coupling two Andreev Qubits via a Cavity Photon



what's next:

- two-qubit gates (timed)
- fidelity, coherence
- address spin degree and try to couple two Andreev **spin qubits**
- can we **improve coherence** (now in the 10 ns range), sweet spots ?

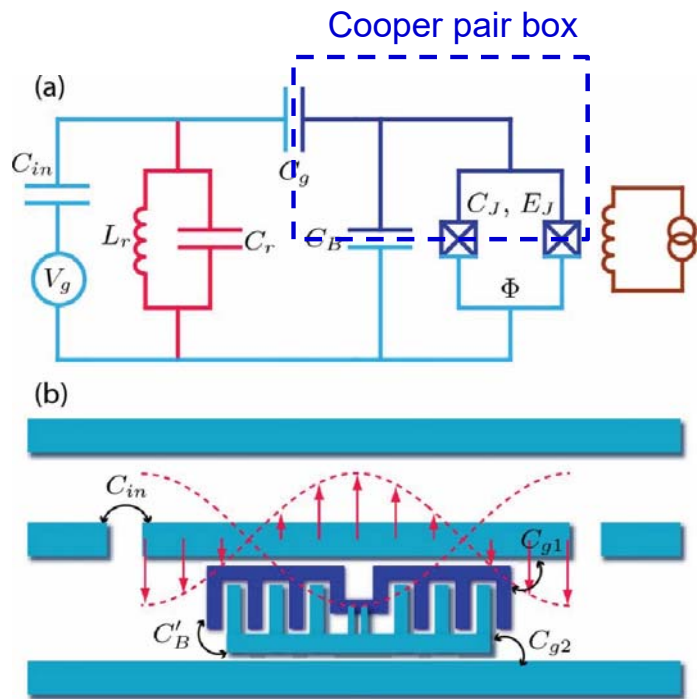
Ge/Si based gatemon as an alternative qubit

Han Zheng, Luk Yi Cheung, Nikunj Sangwan, Tom Jennissen, Artem Kononov, Roy Haller, Joost Ridderbos, Carlo Ciaccia, Jann Hinnerk Ungerer, Erik P.A.M. Bakkers, Andreas Baumgartner, and Christian Schönenberger

EQTC – European Quantum Technology Conference
Oct. 16.10-20.10.2023, Hannover, Germany
by Christian Schönenberger

Quantum- and Nanoelectronics group: www.nanoelectronics.ch
Swiss Nanoscience Institute: <https://nanoscience.ch/en/>
Physics Department of the University of Basel: <https://physik.unibas.ch>

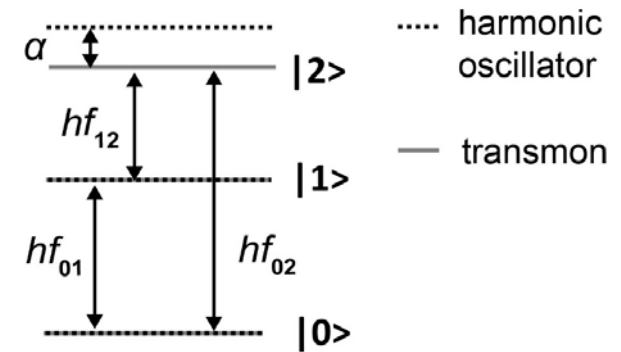
The Transmon Qubit



J. Koch et al. PRA 76, 042319 (2007)

“weakly” non-linear LC circuit
shunted by a large capacitance such that
 $E_J \gg E_c$

$$hf_q = hf_{01} \\ = \sqrt{8E_c E_J} - E_c$$

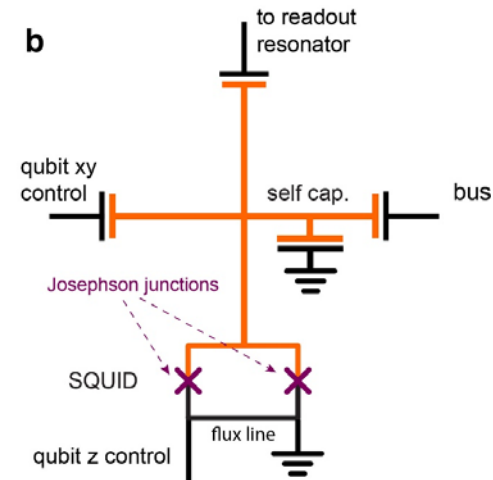
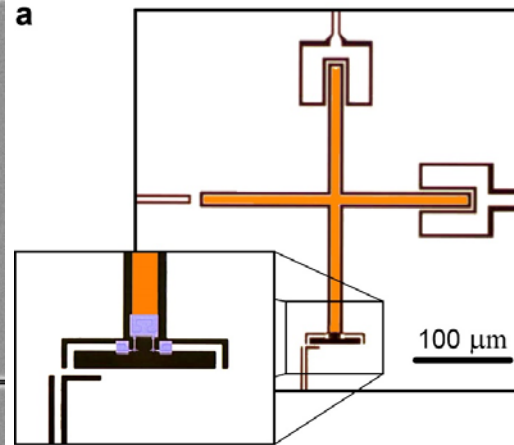
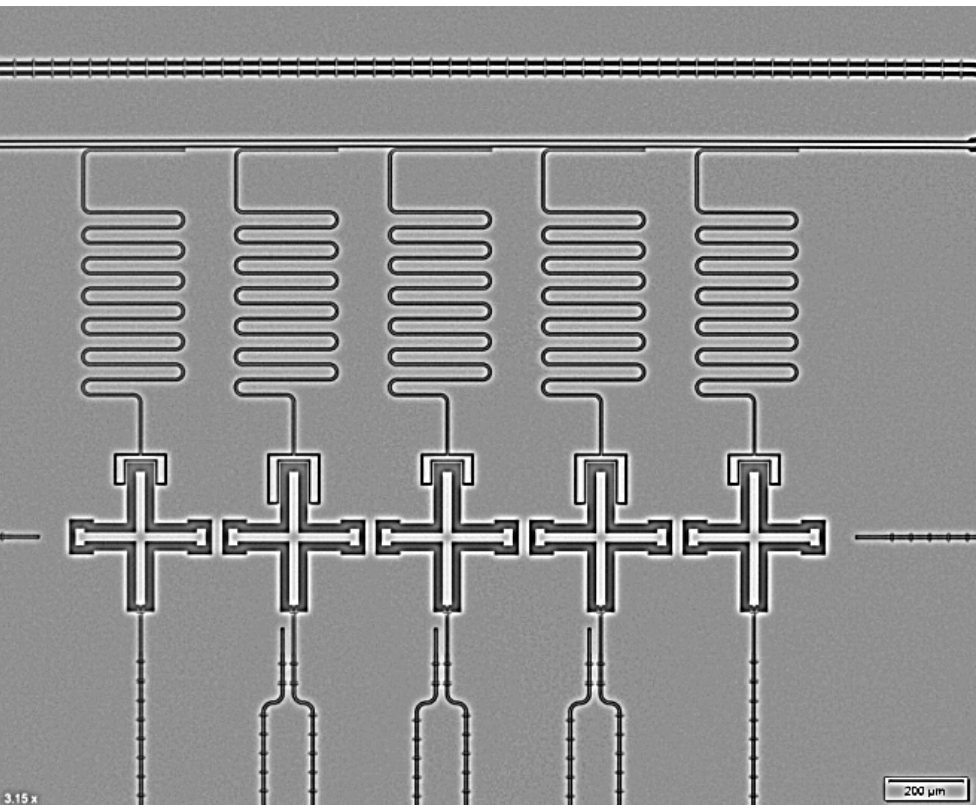


anharmonicity $\alpha := hf_{12} - hf_{01} \cong -E_c$ “small”

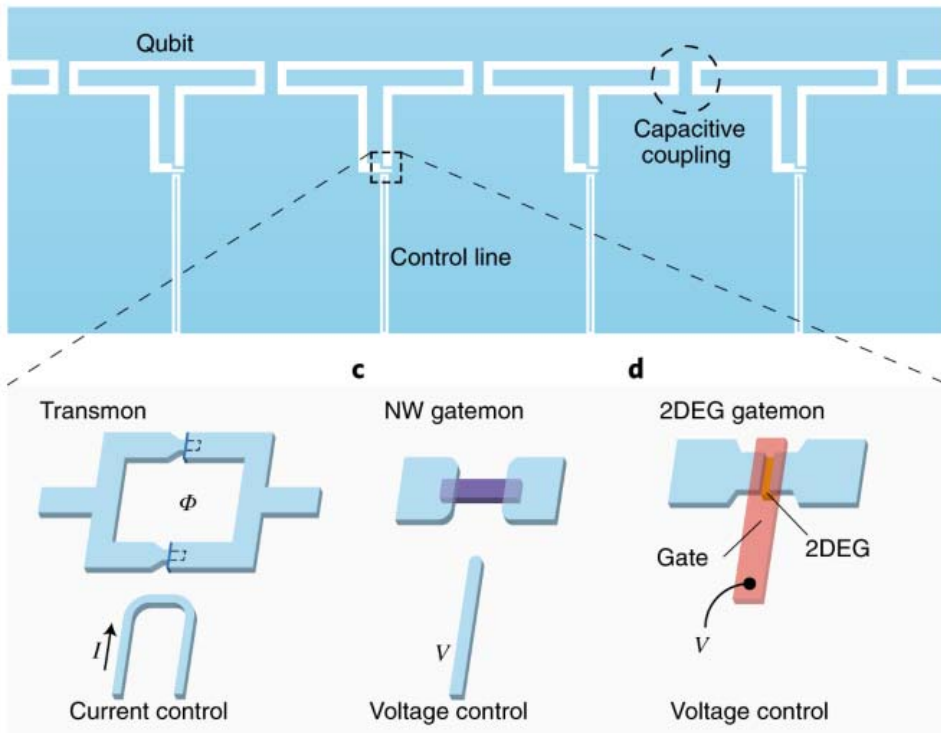
The Transmon Qubit

The **Xmon** qubit

adapted from PRL 111, 080502 (2013) J. Martins' group



Transmon versus Gatemon Qubit



Transmon:

- flux control
- heating (dissipation)
- crosstalk
- small anharmonicity
- “insensitive” to charge noise by construction

Gatemon:

- gate control
- no bias current (lower dissipation)
- voltage is easier to screen
- (even) smaller anharmonicity
- sensitive to charge noise (through gate)

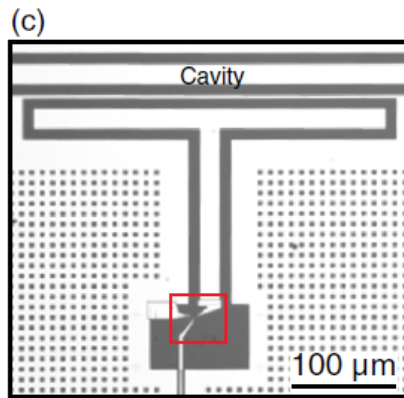
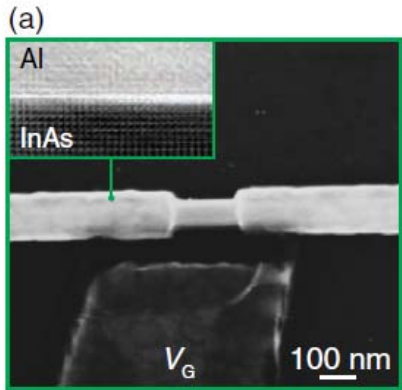
Previous Gatemon Devices

first gatemon devices, both based on InAs nanowires

Semiconductor-Nanowire-Based Superconducting Qubit

2015

T. W. Larsen,¹ K. D. Petersson,¹ F. Kuemmeth,¹ T. S. Jespersen,¹ P. Krogstrup,¹ J. Nygård,^{1,2} and C. M. Marcus¹
¹Center for Quantum Devices, Niels Bohr Institute, University of Copenhagen, Copenhagen 2100, Denmark
²Nano-Science Center, Niels Bohr Institute, University of Copenhagen, Copenhagen 2100, Denmark
 (Received 28 March 2015; published 14 September 2015)



superconductor = Al

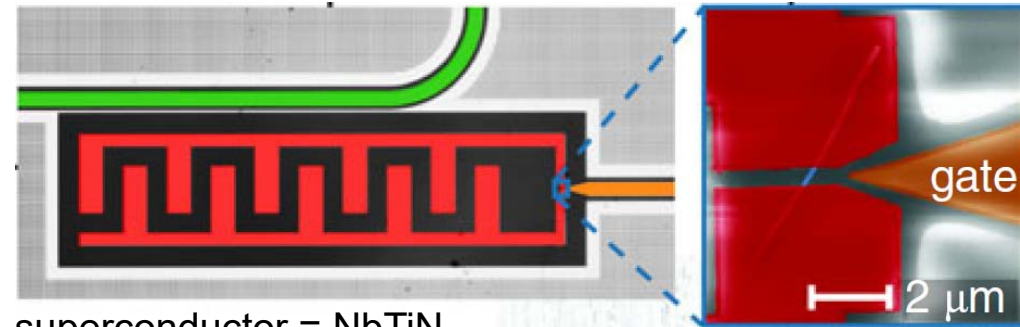
$$T_1 \approx 0.8 \mu s$$

$$T_2^* \approx 1 \mu s$$

Realization of Microwave Quantum Circuits Using Hybrid Superconducting-Semiconducting Nanowire Josephson Elements

2015

G. de Lange,¹ B. van Heck,² A. Bruno,¹ D. J. van Woerkom,¹ A. Geresdi,¹ S. R. Plissard,³
 E. P. A. M. Bakkers,^{1,3} A. R. Akhmerov,¹ and L. DiCarlo¹
¹QuTech and Kavli Institute of Nanoscience, Delft University of Technology, 2600 GA Delft, The Netherlands
²Instituut-Lorentz, Leiden University, 2300 RA Leiden, The Netherlands
³Department of Applied Physics, Eindhoven University of Technology, 5600 MB Eindhoven, The Netherlands
 (Received 30 March 2015; published 14 September 2015)



superconductor = NbTiN

looked additionally into SQUID junctions tuned to half-flux quantum (different spectrum, fluxonium)

Previous Gatemon Devices

first gatemon devices, both based on InAs nanowires

PRL 116, 150505 (2016) PHYSICAL REVIEW LETTERS week ending 15 APRIL 2016

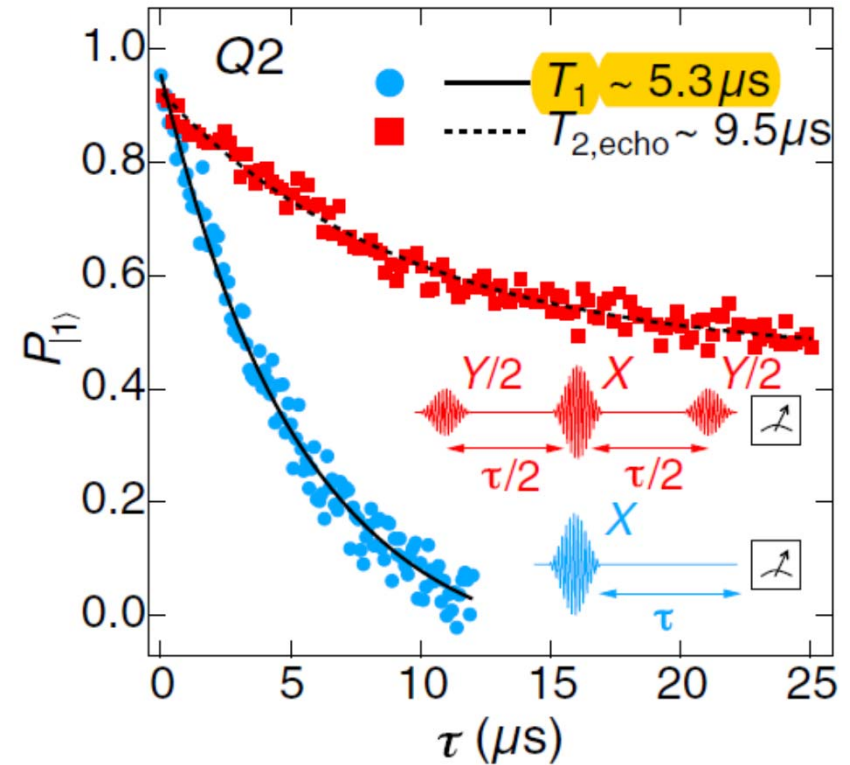
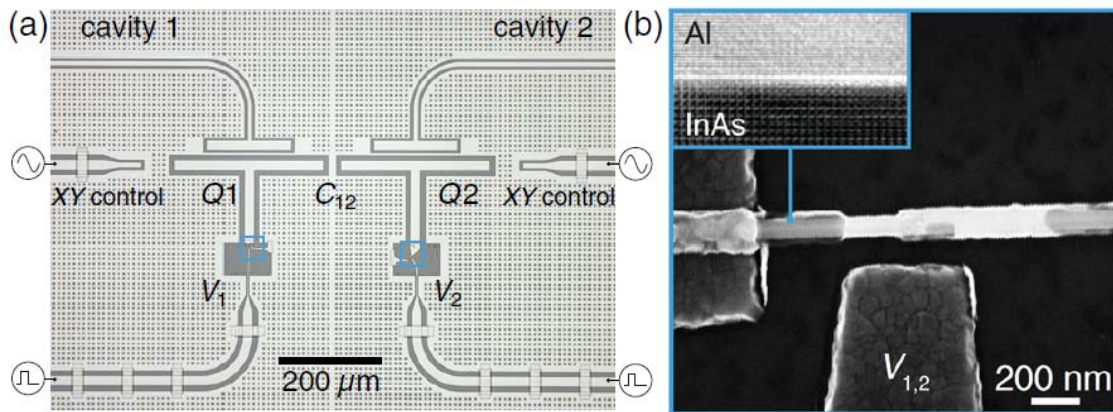
Gatemon Benchmarking and Two-Qubit Operations

2016

L. Casparis,¹ T. W. Larsen,¹ M. S. Olsen,¹ F. Kuemmeth,¹ P. Krogstrup,¹ J. Nygård,^{1,2}
K. D. Petersson,¹ and C. M. Marcus¹

¹Center for Quantum Devices, Station Q Copenhagen, Niels Bohr Institute, University of Copenhagen, Copenhagen DK-2100, Denmark

²Nano-Science Center, Niels Bohr Institute, University of Copenhagen, Copenhagen DK-2100, Denmark
(Received 30 December 2015; published 15 April 2016)



Previous Gatemon Devices

first gatemon devices, both based on InAs nanowires

PHYSICAL REVIEW LETTERS **120**, 100502 (2018)

Evolution of Nanowire Transmon Qubits and Their Coherence in

F. Luthi,^{1,2} T. Stavenga,^{1,2} O. W. Enzing,^{1,2} A. Bruno,^{1,2} C. Dickel,^{1,2} N. K. Lang,¹
T. S. Jespersen,³ J. Nygård,^{3,4} P. Krogstrup,³ and L. DiCarlo¹

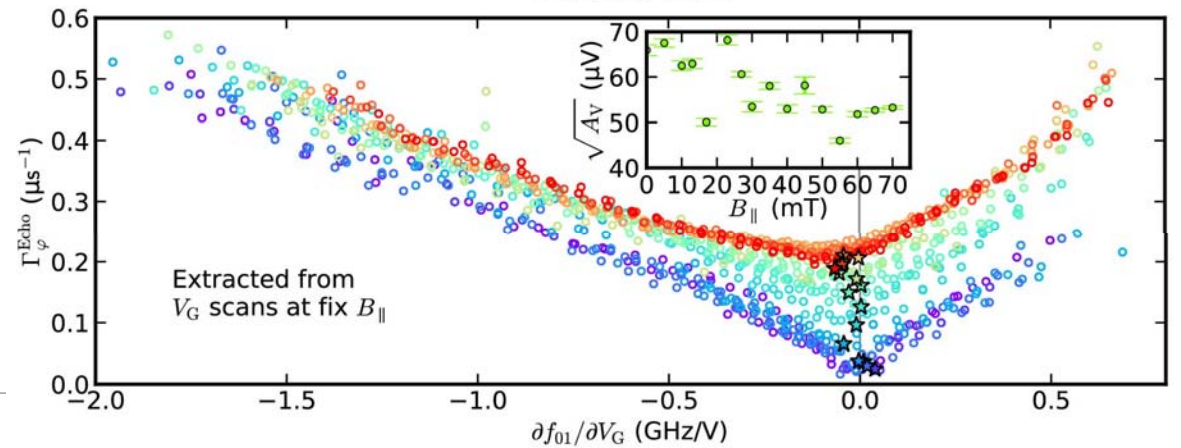
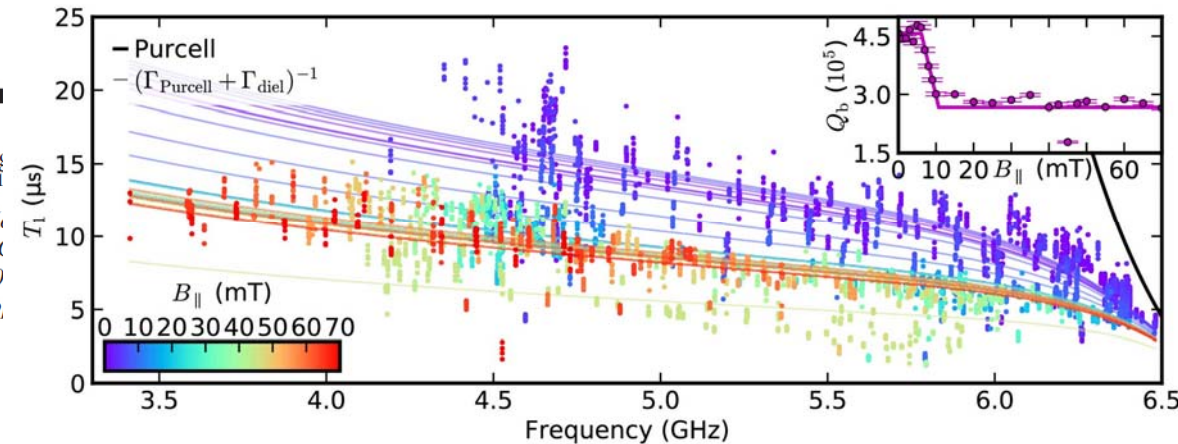
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 (Received 22 November 2017; published 9 March 2018)



obtained $T_2^{echo} \cong 40 \mu s$
at sweet spot and at $B = 0!$

Previous Gatemon Devices

InAs quantum wells

nature
nanotechnology

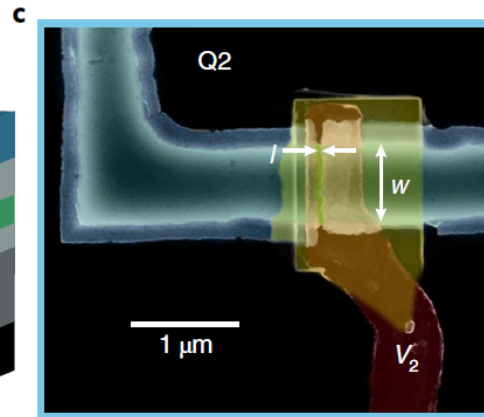
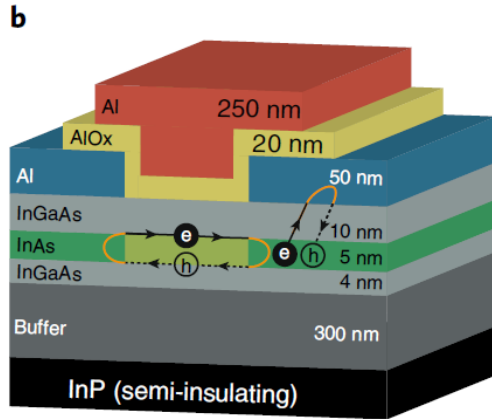
LETTERS

<https://doi.org/10.1038/s41565-018-0207-y>

2018

Superconducting gatemon qubit based on a proximitized two-dimensional electron gas

Lucas Casparis^{1,8}, Malcolm R. Connolly^{1,8}, Morten Kjaergaard^{1,7}, Natalie J. Pearson^{1,2}, Anders Kringhøj¹, Thorvald W. Larsen¹, Ferdinand Kuemmeth¹, Tiantian Wang^{3,4}, Candice Thomas^{3,4}, Sergei Gronin⁴, Geoffrey C. Gardner⁴, Michael J. Manfra^{3,4,5,6}, Charles M. Marcus¹ and Karl D. Petersson^{1*}



$$T_1 \cong 1 \mu\text{s} \quad T_2^* \cong 400 \text{ ns} \quad T_2^{\text{echo}} \cong 2.2 \mu\text{s}$$

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Graphene Gatemon

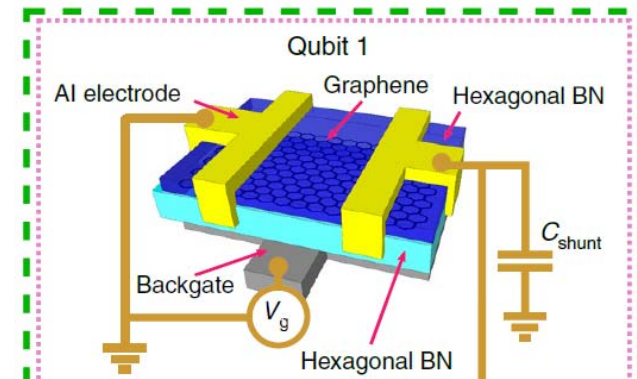
LETTERS

<https://doi.org/10.1038/s41565-018-0329-2>

2019

Coherent control of a hybrid superconducting circuit made with graphene-based van der Waals heterostructures

Joel I-Jan Wang^{1,7*}, Daniel Rodan-Legrain^{2,7}, Landry Bretheau³, Daniel L. Campbell¹, Bharath Kannan^{1,4}, David Kim⁵, Morten Kjaergaard¹, Philip Krantz¹, Gabriel O. Samach^{4,5}, Fei Yan¹, Jonilyn L. Yoder⁵, Kenji Watanabe⁶, Takashi Taniguchi⁶, Terry P. Orlando^{1,4}, Simon Gustavsson¹, Pablo Jarillo-Herrero^{2*} and William D. Oliver^{1,2,5*}



$$T_1 \cong 36 \text{ ns}$$

$$T_2^* \cong 55 \text{ ns}$$

Previous Gatemon Devices

Carbon Nanotube Gatemon

2021

PHYSICAL REVIEW APPLIED 15, 064050 (2021)

Circuit Quantum Electrodynamics with Carbon-Nanotube-Based Superconducting Quantum Circuits

Matthias Mergenthaler^{1,2,*}, Ani Nersisyan¹, Andrew Patterson¹, Martina Esposito¹,
Andreas Baumgartner³, Christian Schönenberger³, G. Andrew D. Briggs², Edward A. Laird^{4,2} and
Peter J. Leek^{1,†}

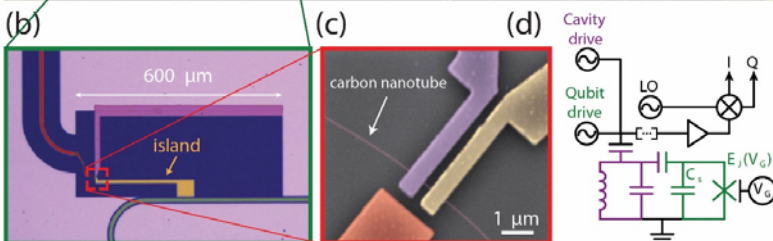
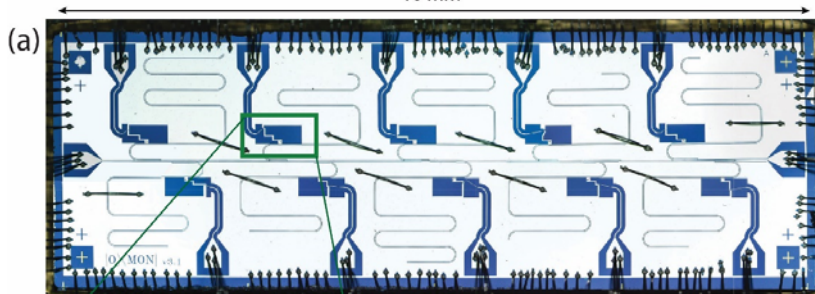
¹Clarendon Laboratory, Department of Physics, University of Oxford, Oxford OX1 3PU, United Kingdom

²Department of Materials, University of Oxford, Oxford OX1 3PH, United Kingdom

³Department of Physics, University of Basel, Klingelbergstrasse 82, Basel CH-4056, Switzerland

⁴Department of Physics, Lancaster University, Lancaster LA1 4YB, United Kingdom

10 mm



$$T_1 < 200 \text{ ns}$$

$$T_2' \cong 10 \text{ ns}$$

(linewidth)

Two examples of type-IV materials (graphene and CNT) for hybrid super-semi Josephson junctions (JJs).

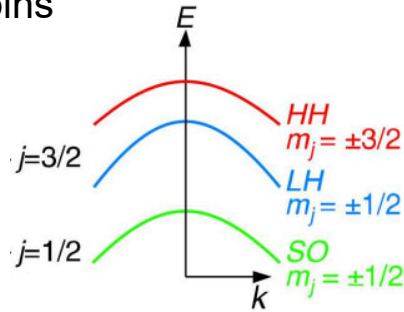
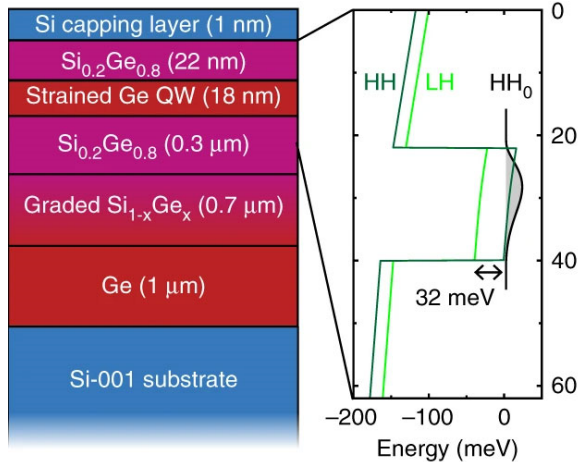
What about **Silicon** and **Germanium**?

Type-IV semiconductors

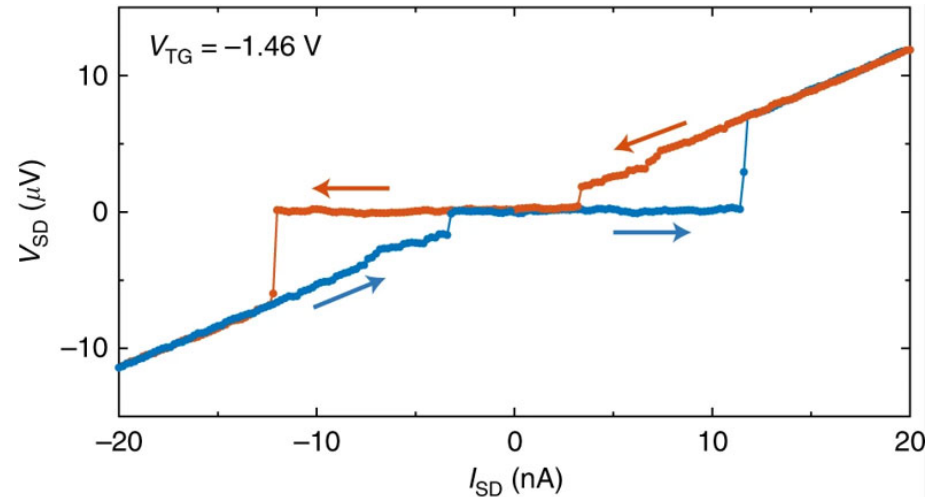
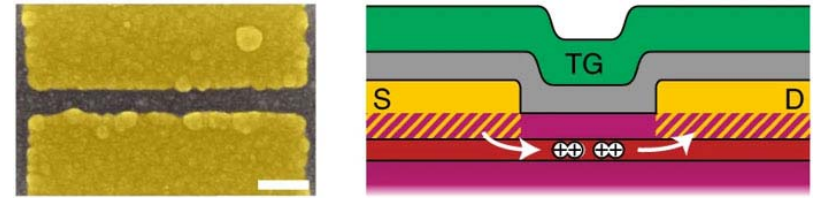
Scappucci and Veldhorst et al.
Gate-controlled quantum dots and superconductivity in planar germanium: Nature Comm. 9;2835 (2018)

Germanium:

- highest hole mobility $> 1 \cdot 10^6 \text{ cm}^2/\text{Vs}$
- no valley structure (a single band maximum)
- strained
- gate-tunable large spin-orbit interaction
- very low coupling to nuclear spins



induced superconductivity through Al contacts



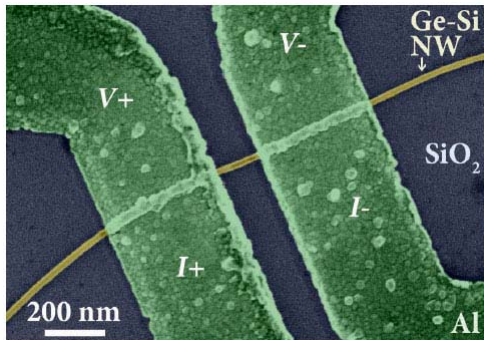
but, in this early results the $I_C \cdot R_N$ product was way lower than the ideal ballistic limit of:

$$I_C = e\Delta/\hbar \text{ (per channel)} \rightarrow I_C \cdot R_N = \pi\Delta/e$$

Germanium Josephson FET

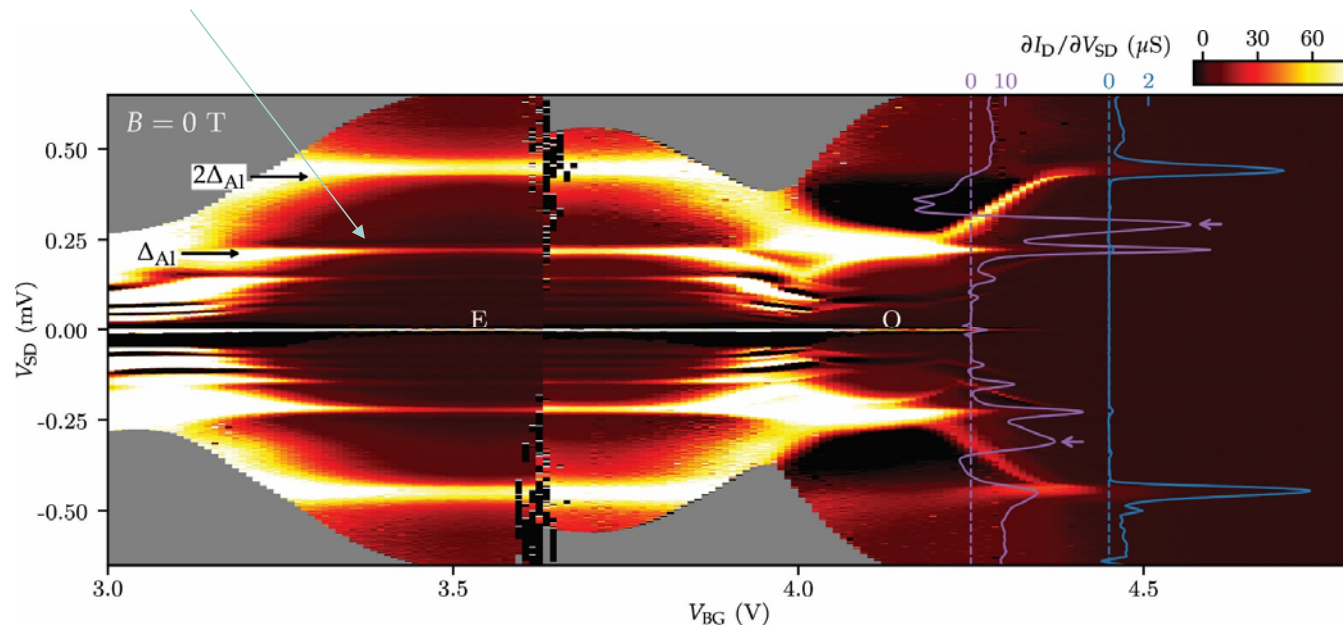
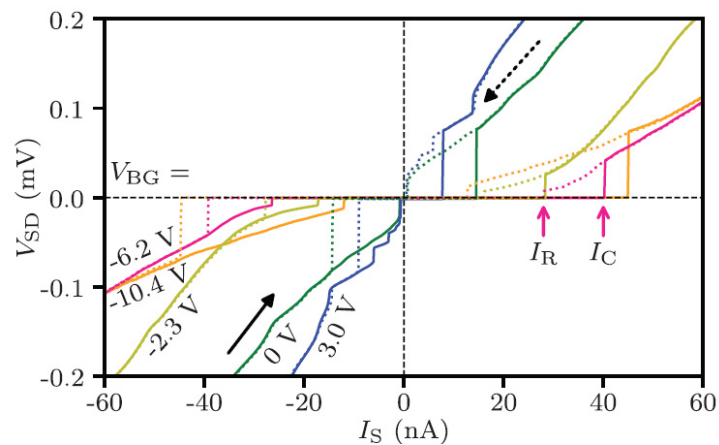
Josephson Effect in a Few-Hole Quantum Dot

Joost Ridderbos et al. Adv. Mat. 30: 1802257 (2018) & Nano Lett. 20: 122 (2020)

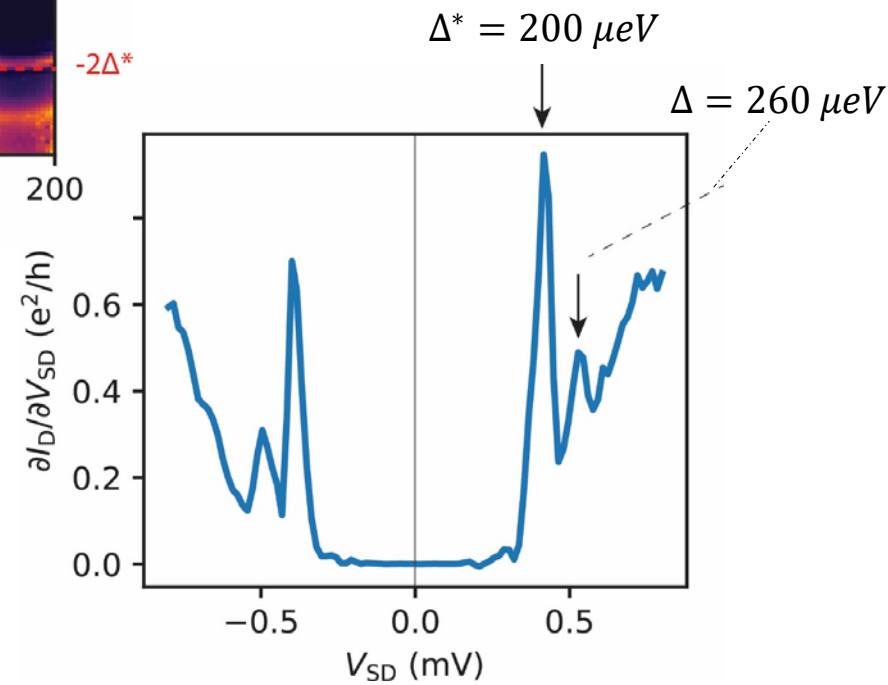
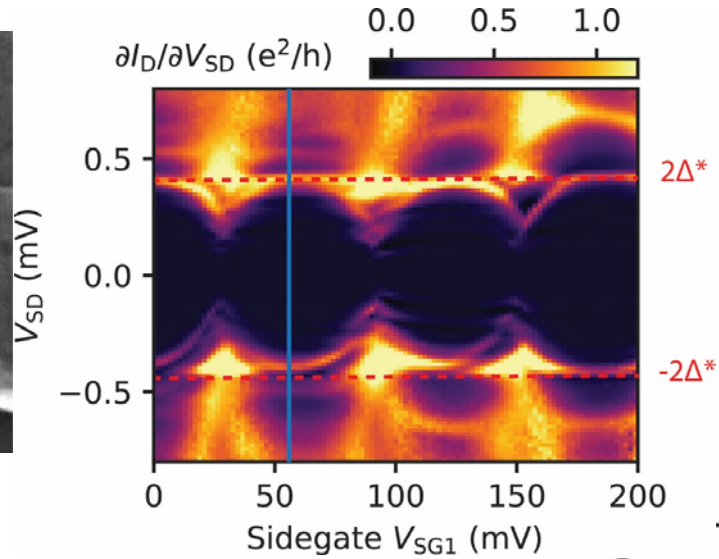
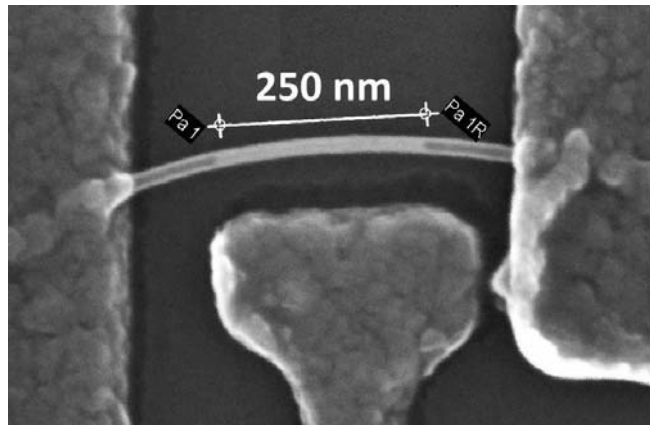


GeSi core-shell nanowires (from TUE, E. Bakker's group)

- large $I_c \cdot R_N$ product of order $200 \mu eV$
- hard gap in tunnelling regime
- **multiple Andreev reflection (MAR)** at larger transmission probability

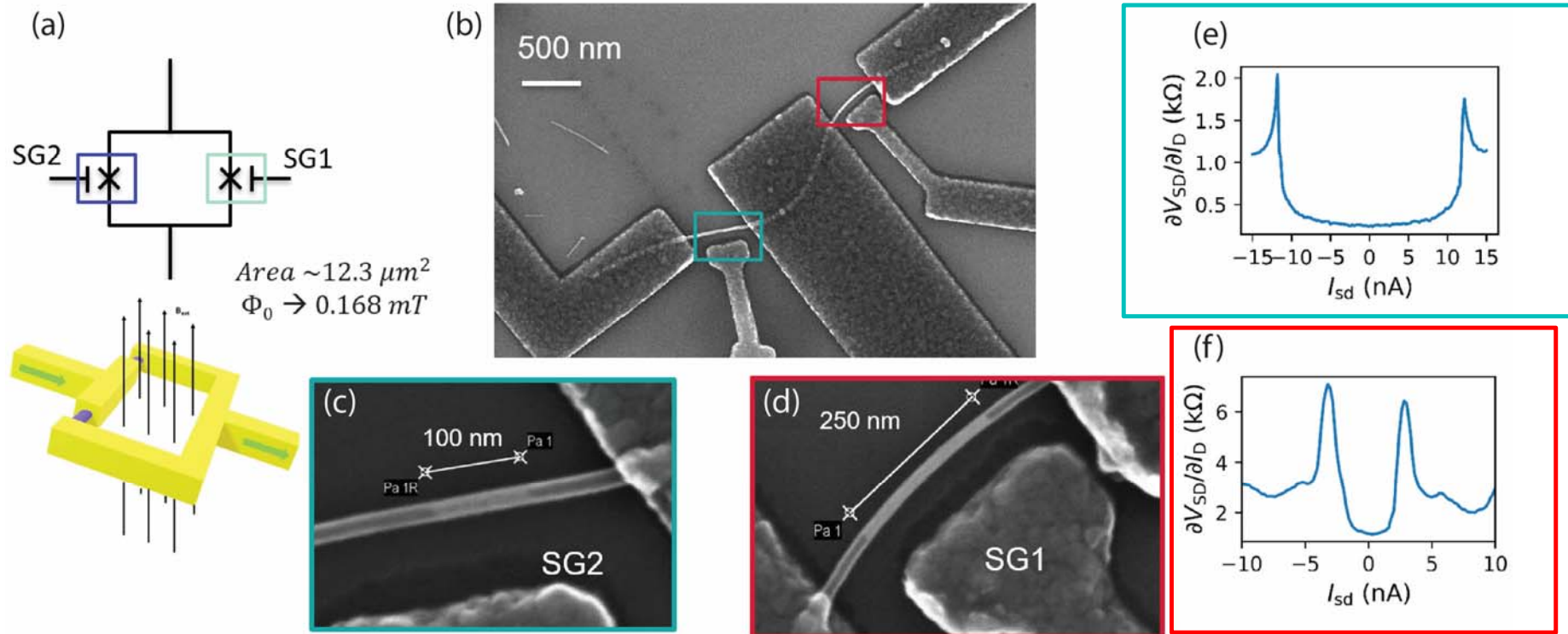


Germanium Josephson FET



Han Zheng, Tom Jennissen, J. Ridderbos (unpublished)

Germanium Josephson FET

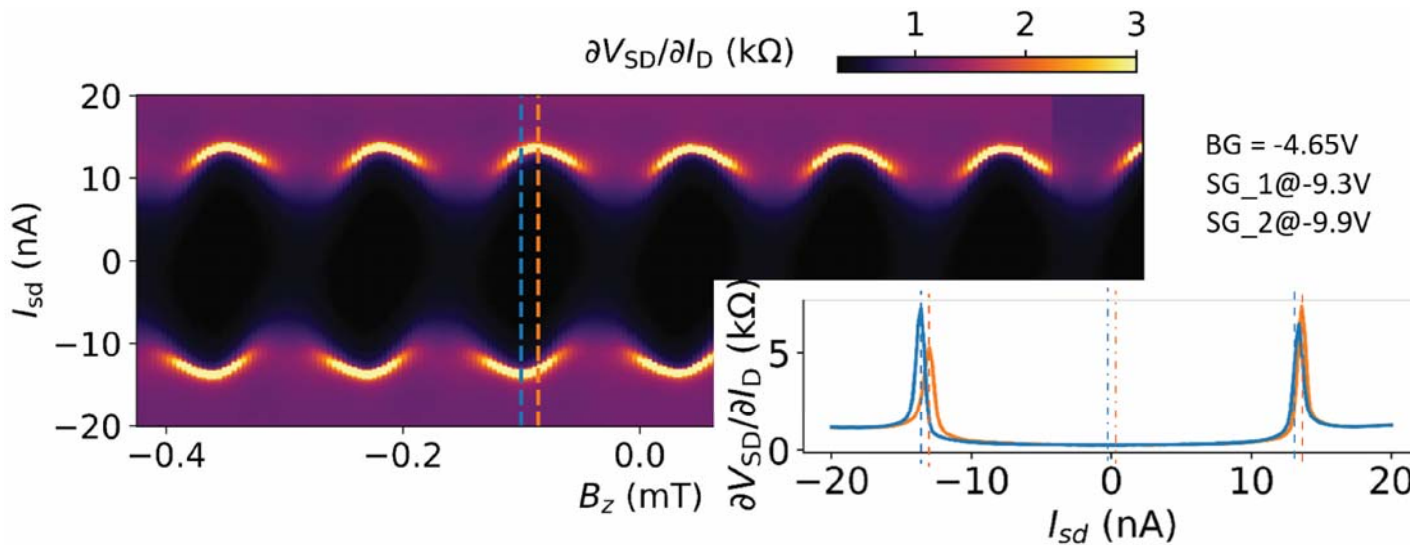


Han Zheng, Tom Jennissen, J. Ridderbos (unpublished)

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Current-Phase Relation of a GeSi core-shell nanowire JJ

In an asymmetric SQUID with a strong reference JJ having at least 10 times higher critical current, one can obtain the CPR of the weaker junction by measuring the critical current of the SQUID as a function of flux through the SQUID loop.



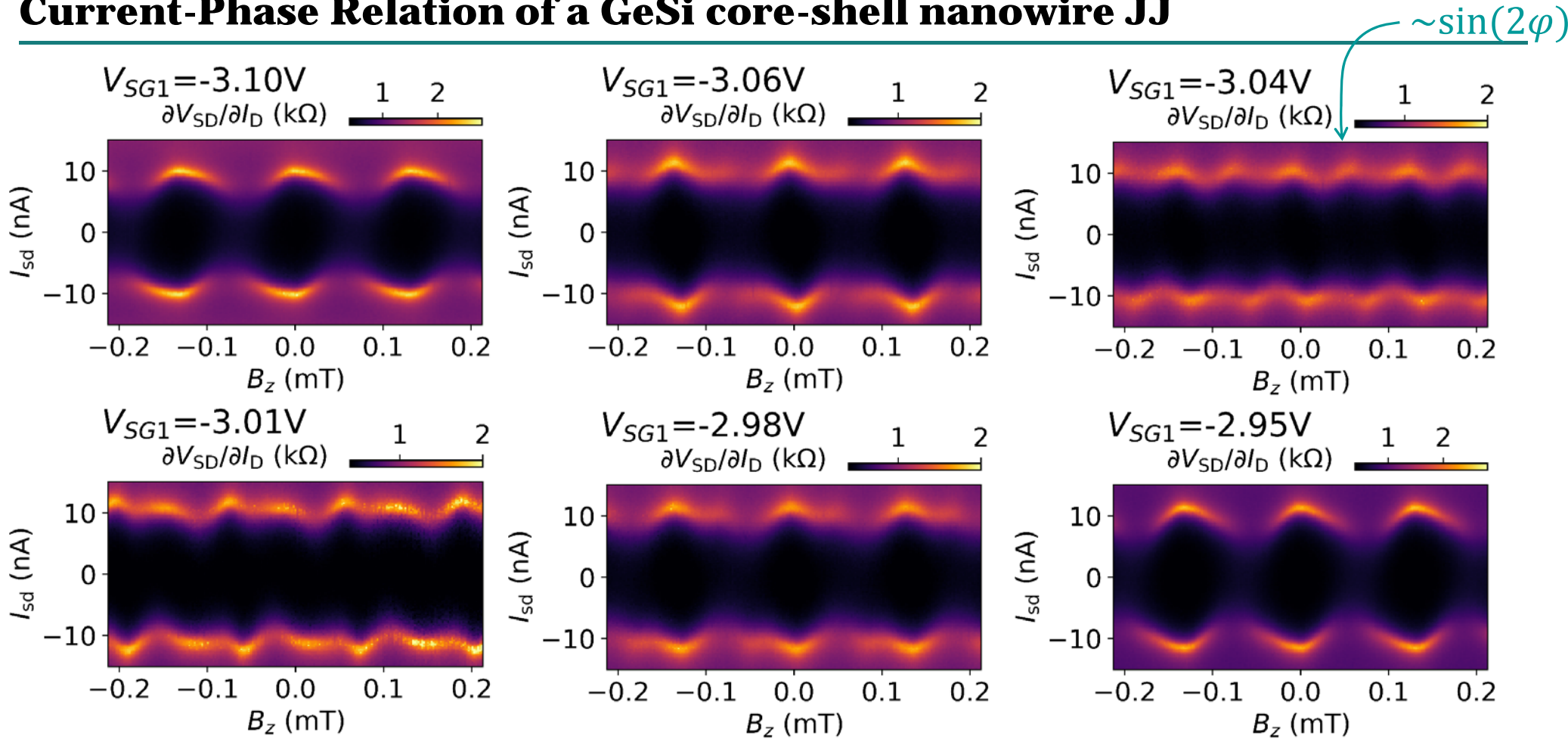
observe:

- non-sinusoidal CPR (indication of transparent JJ)
- diode effect (also caused by non-sinusoidal character)
- good tunability of supercurrent

Diode effect due to non-sinusoidal CPR. See: C. Ciaccia et al. (Schönenberger group), Phys. Rev. Research 5, 33131 (2023) and see also M. Valentini et al. (Katsaros' group) arXiv:2306.07109.

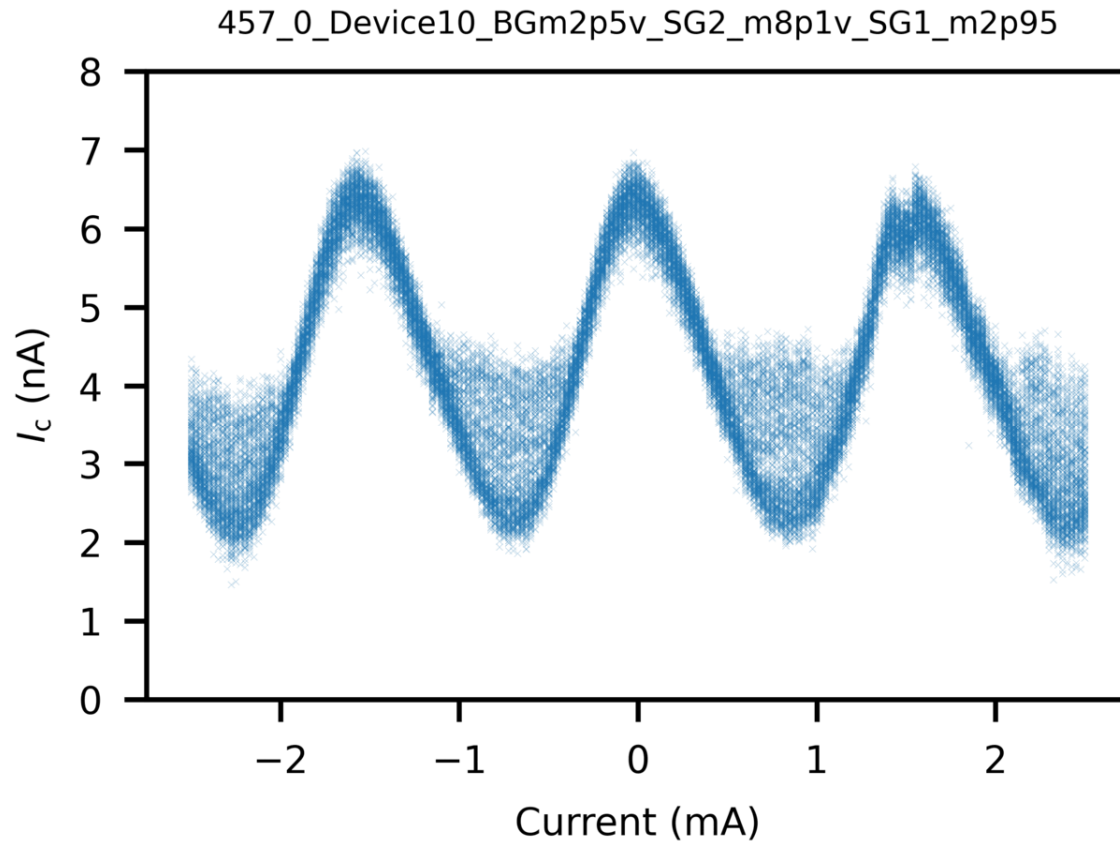
Han Zheng, Tom Jennissen, J. Ridderbos (unpublished)

Current-Phase Relation of a GeSi core-shell nanowire JJ



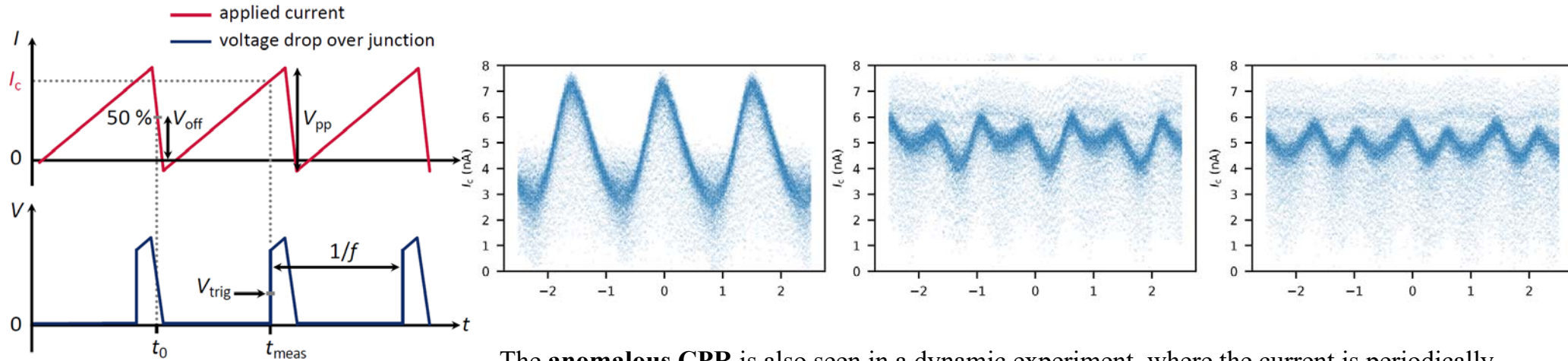
Han Zheng, Tom Jennissen, J. Ridderbos (unpublished)

Current-Phase Relation of a GeSi core-shell nanowire JJ

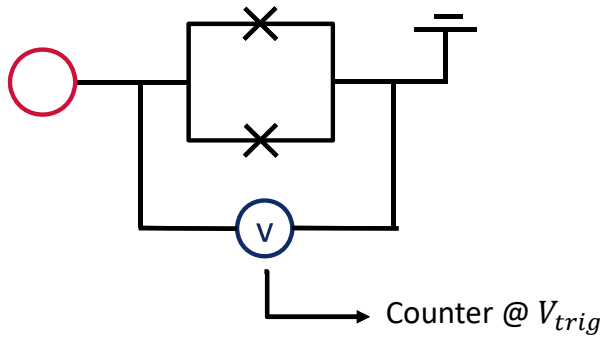


Han Zheng, Tom Jennissen, J. Ridderbos (unpublished)

Current-Phase Relation of a GeSi core-shell nanowire JJ



The **anomalous CPR** is also seen in a dynamic experiment, where the current is periodically swept up and down

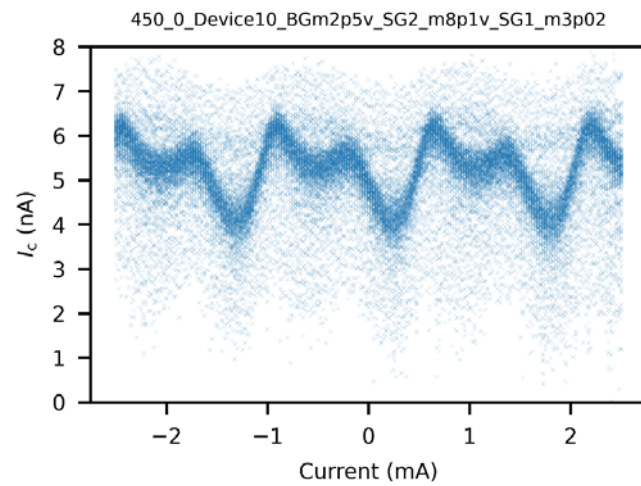
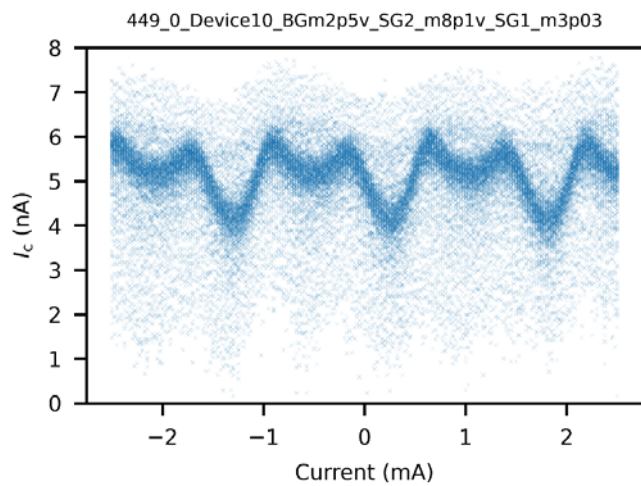
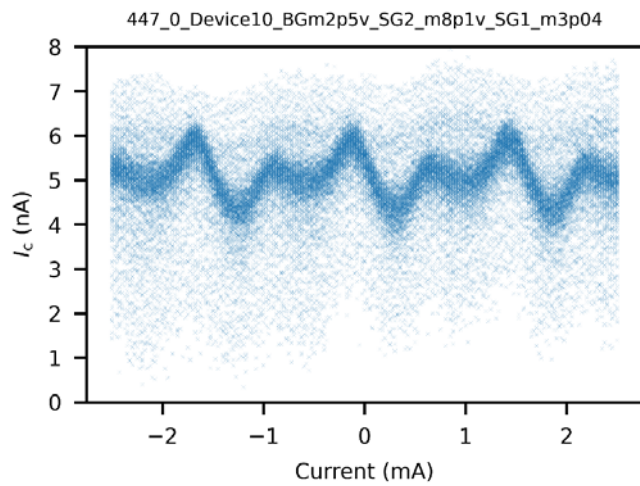
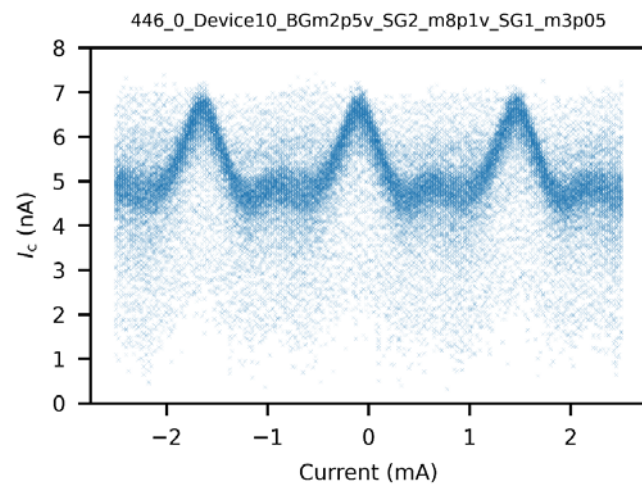
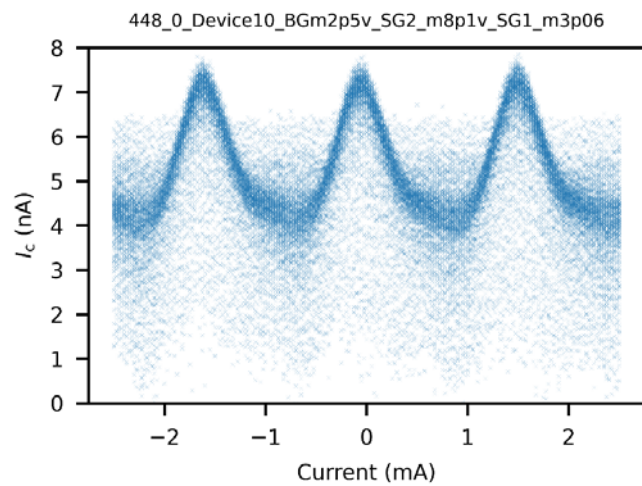
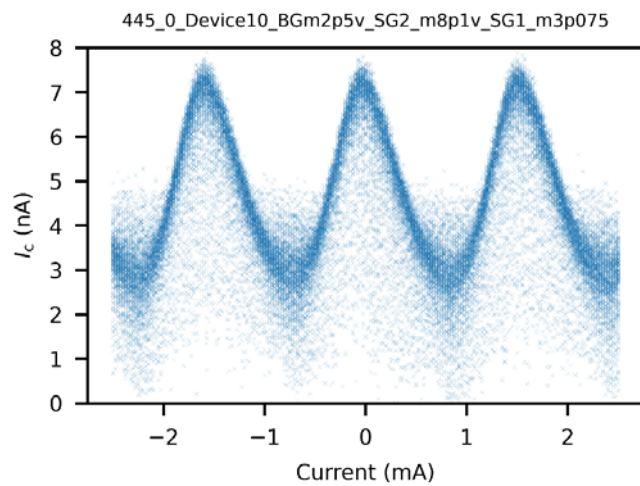


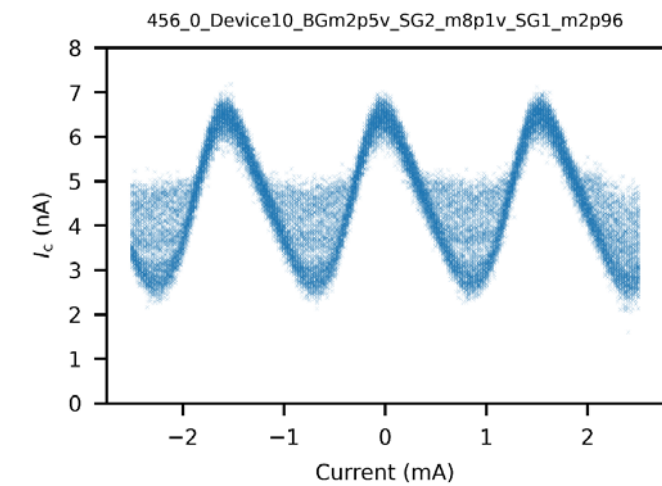
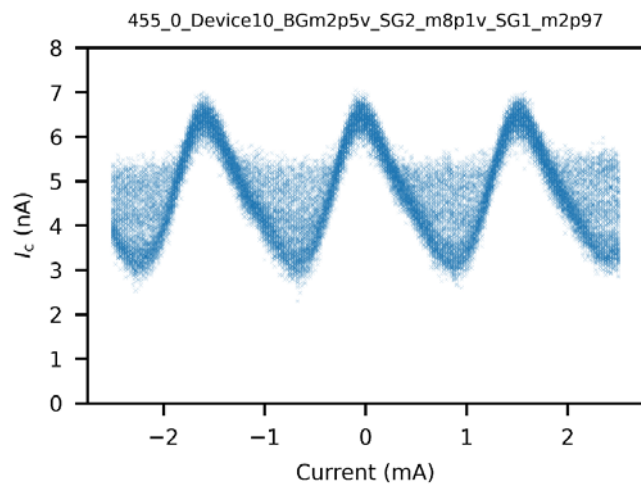
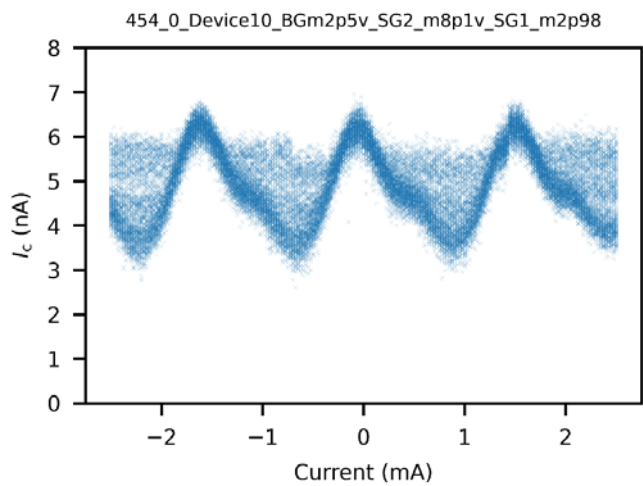
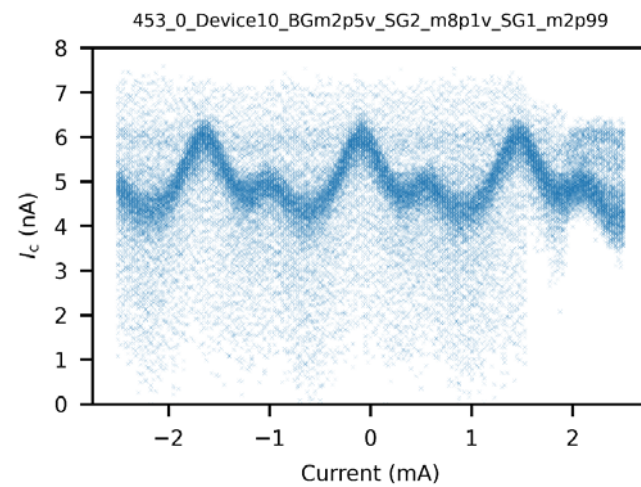
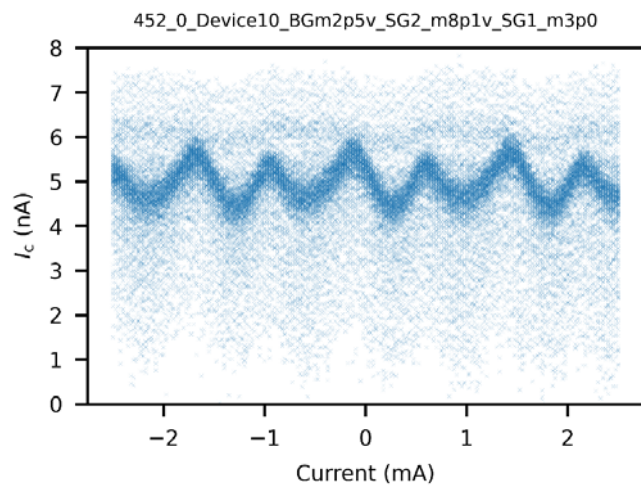
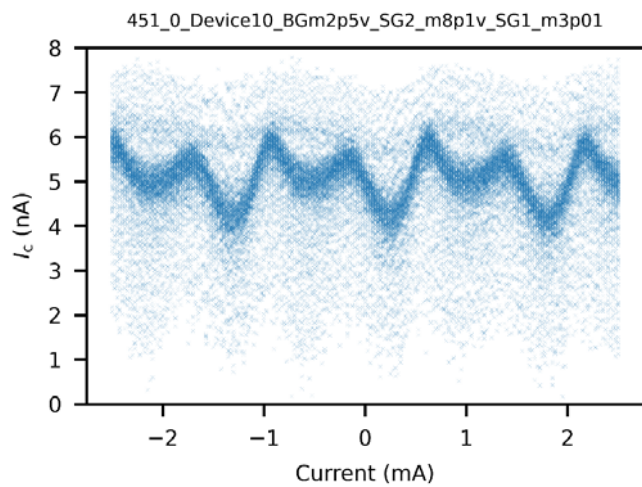
Origin:

1. Spin-orbit effect combined with magnetic field
2. Interference of the supercurrent of two Andreev-bound states for which the fundamental $\sin(\varphi)$ contributions cancel each other (has application for a parity protected qubit)

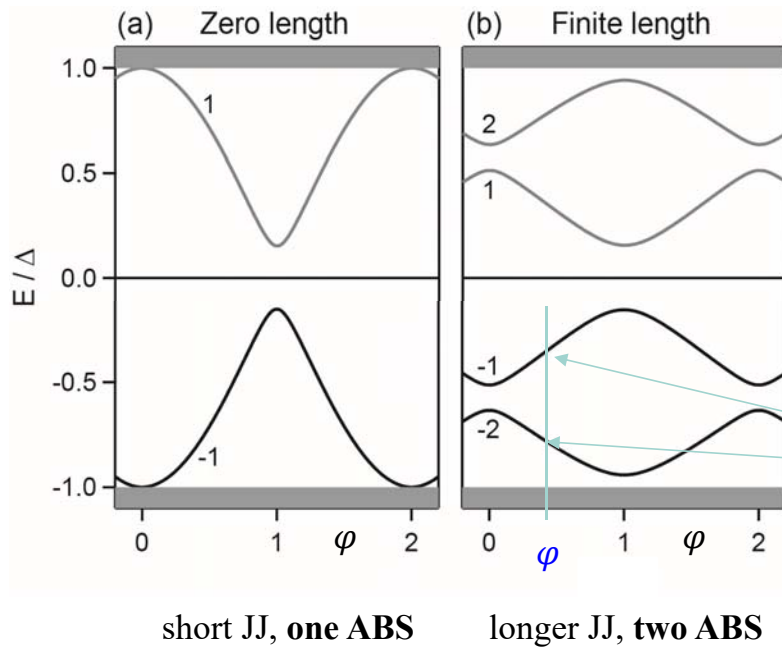
Han Zheng, Tom Jennisen, J. Ridderbos (unpublished)

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Current-Phase Relation of a GeSi core-shell nanowire JJ



~ $\sin(2\varphi)$ junction ?

$$I(\varphi) = \frac{2\pi}{\Phi_0} \frac{\partial E(\varphi)}{\partial \varphi}$$

opposite slope → partial cancelation of supercurrent*

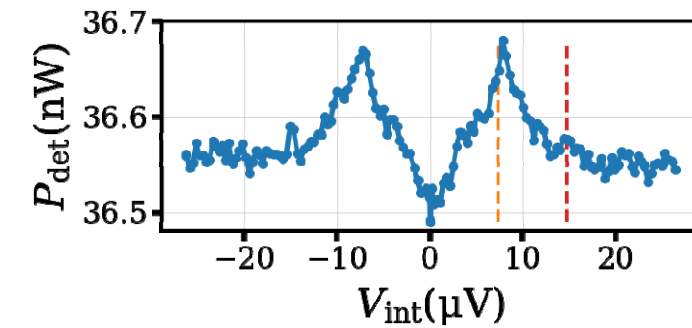
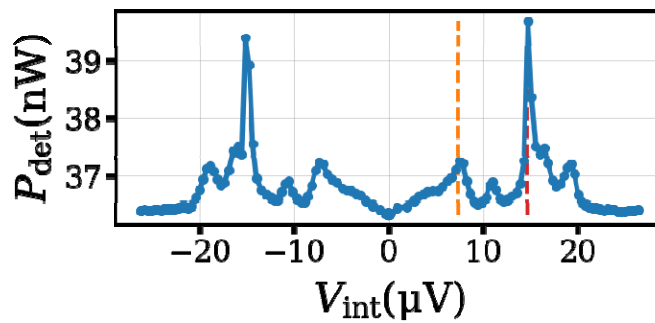
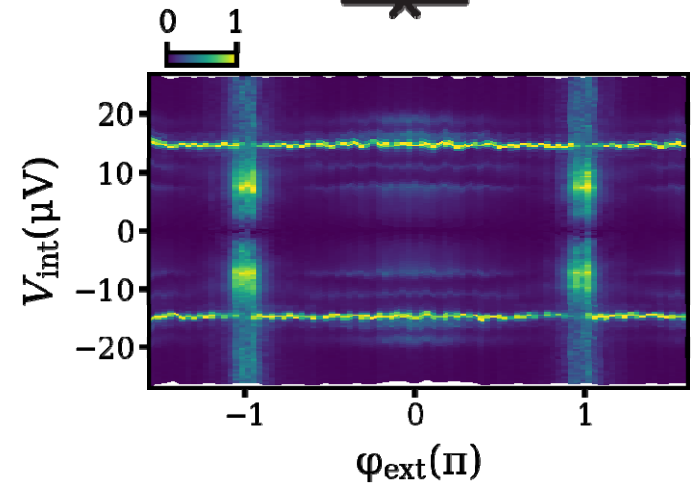
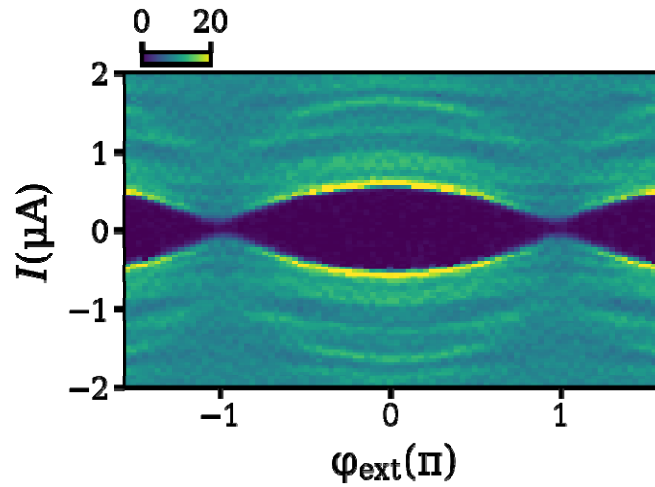
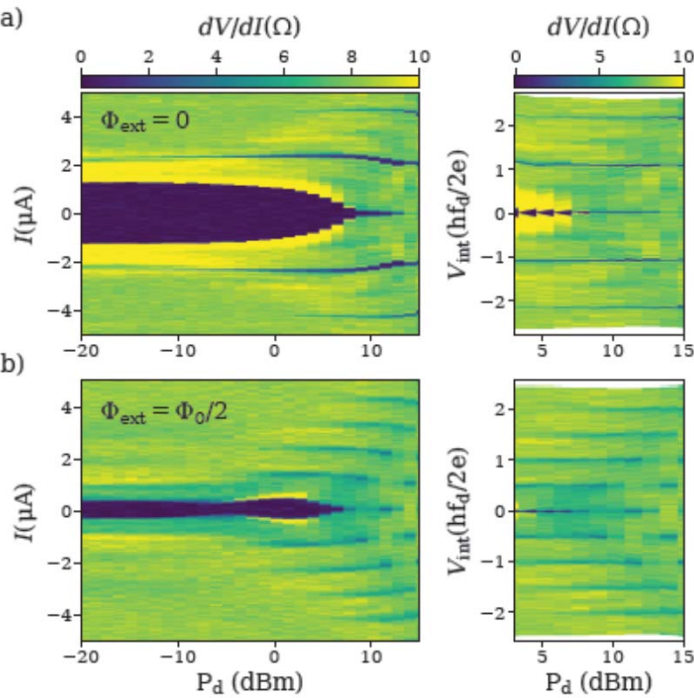
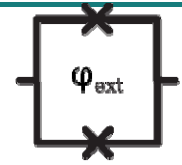
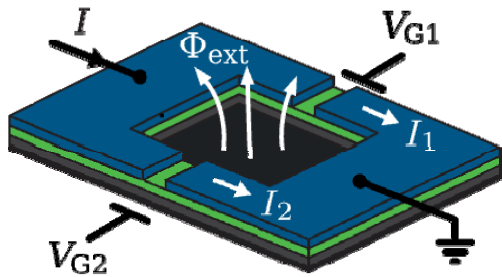
Hence, it can happen that the fundamental contributions $\sin(\varphi)$ of the two ABSs cancel each other. The 1st non-zero harmonics would then follow $\sin(2\varphi)$ → **$\sin(2\varphi)$ junction**

*idea by Jelena Klinovaja

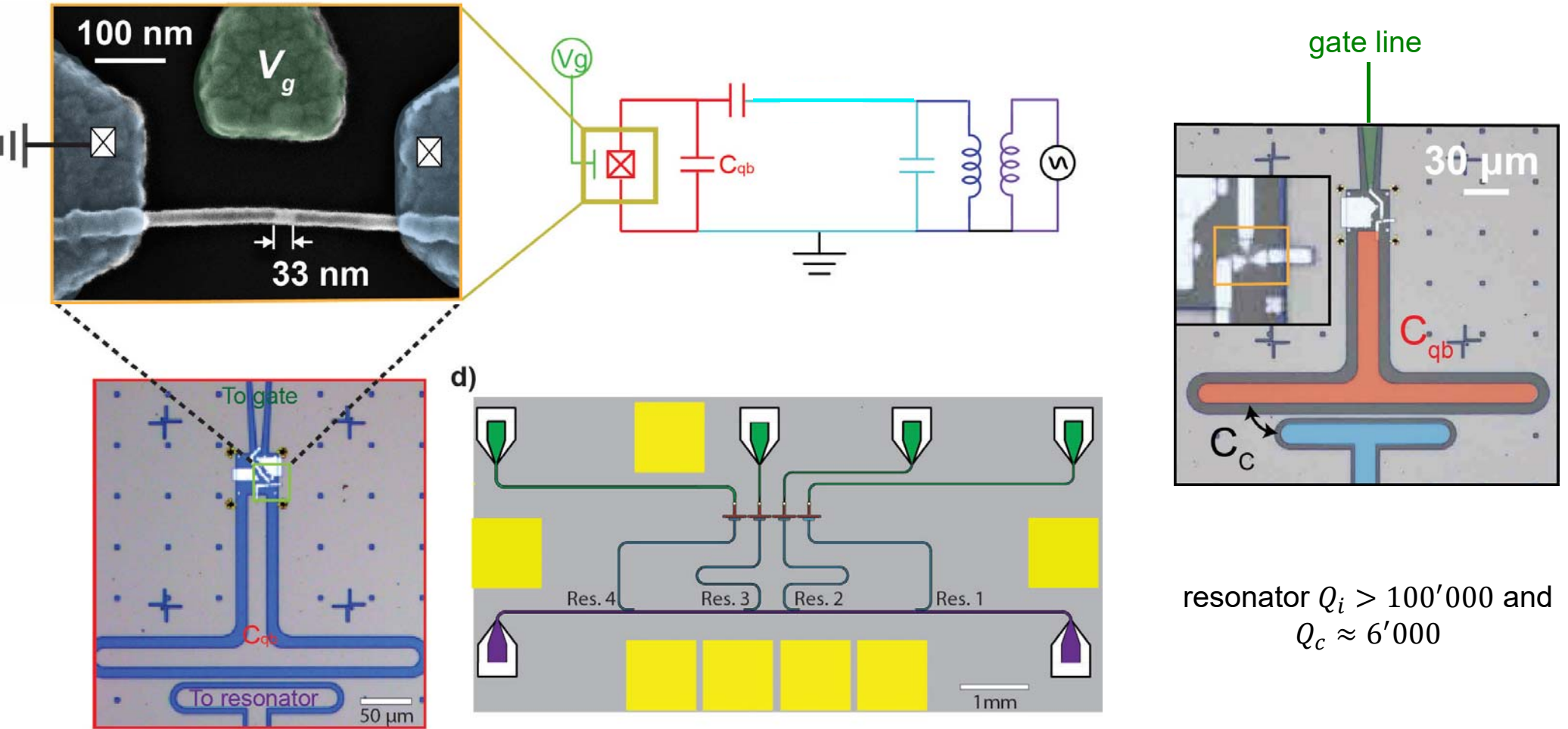
Han Zheng, Tom Jennissen, J. Ridderbos (unpublished)

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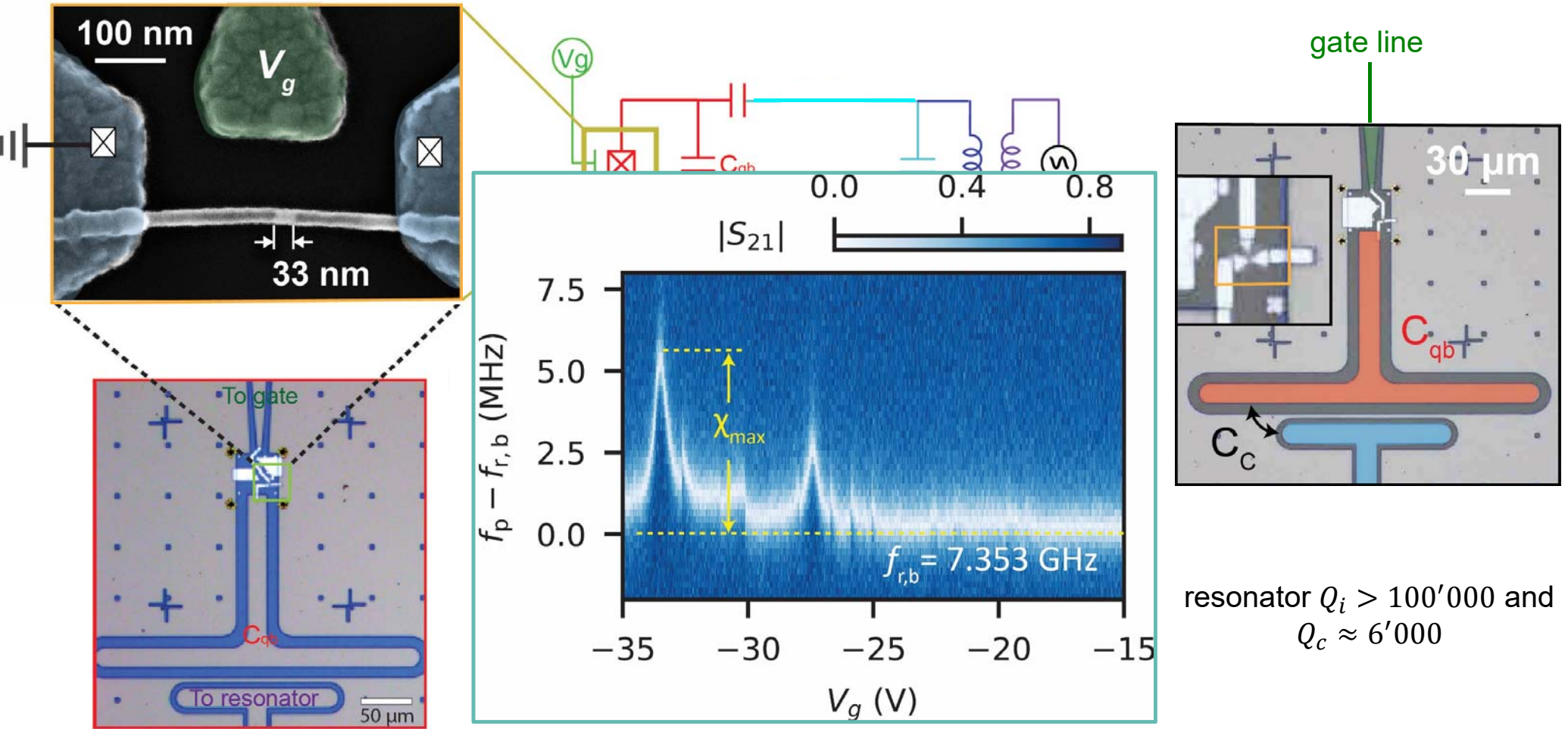
Engineering a $\cos(2\phi)$ junction with a SQUID



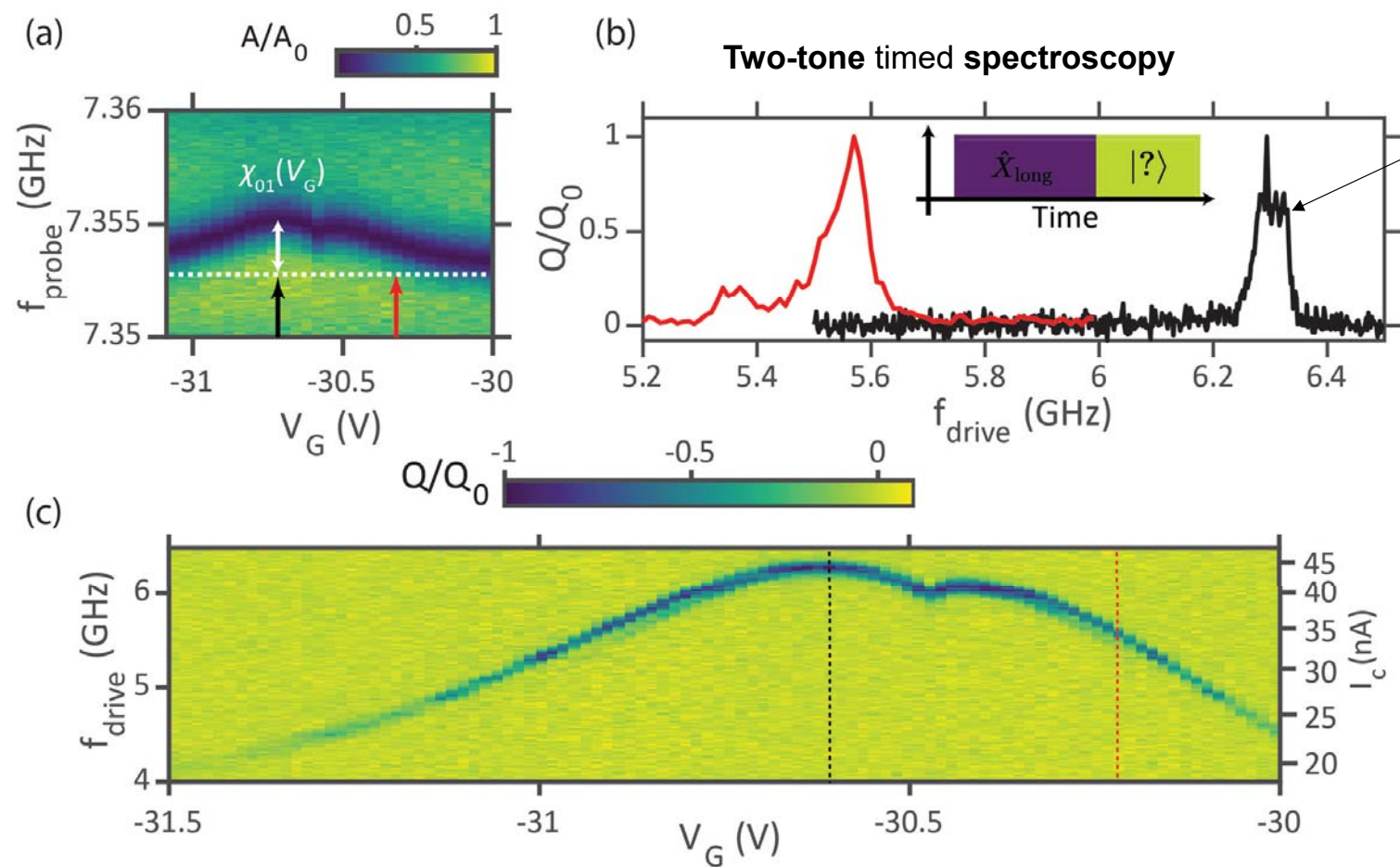
GeSi core-shell nanowire Gatemon



GeSi core-shell nanowire Gatemon



GeSi core-shell nanowire Gatemon

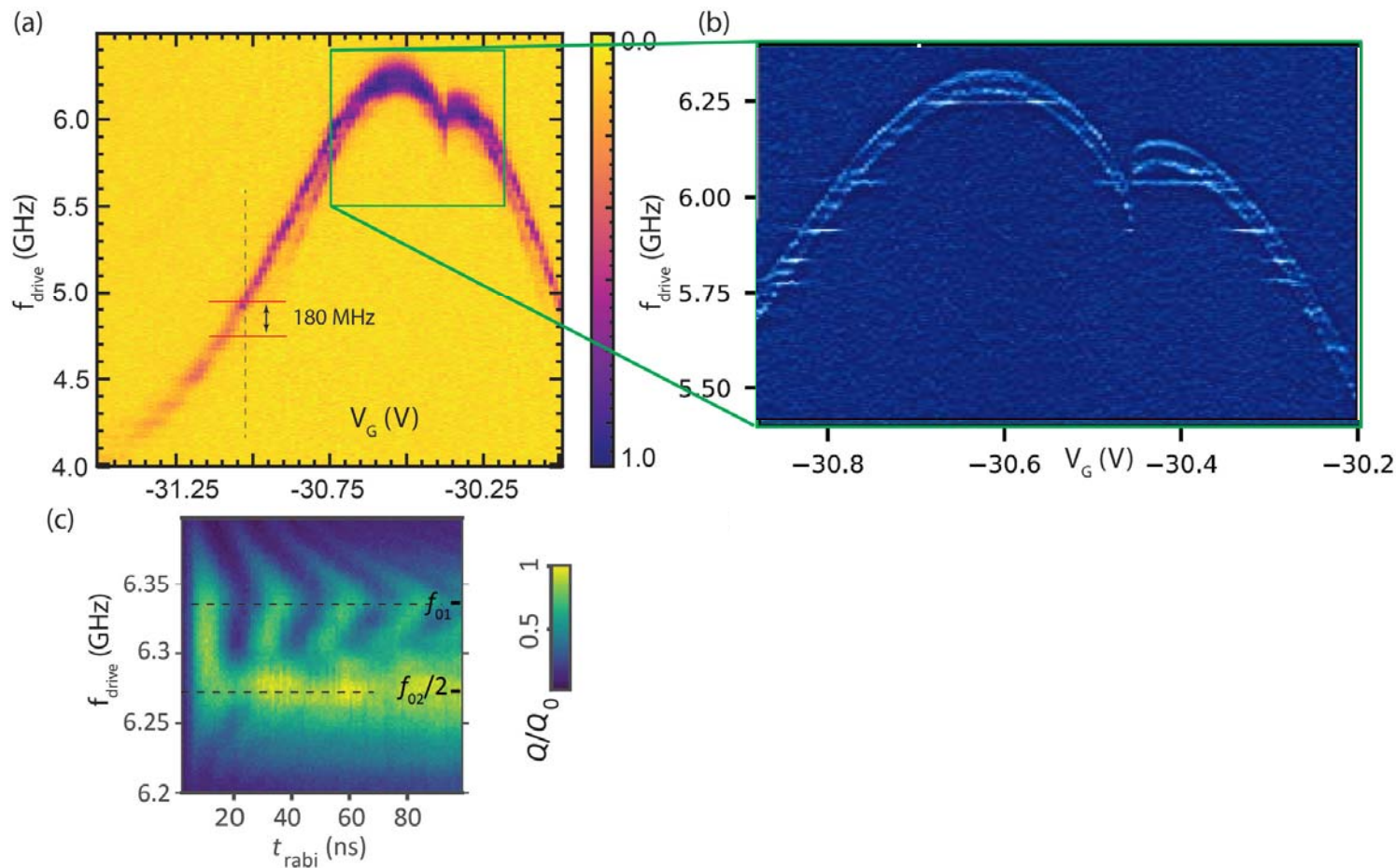


not a single Lorentzian

from dispersive shift χ and detuning we obtain:

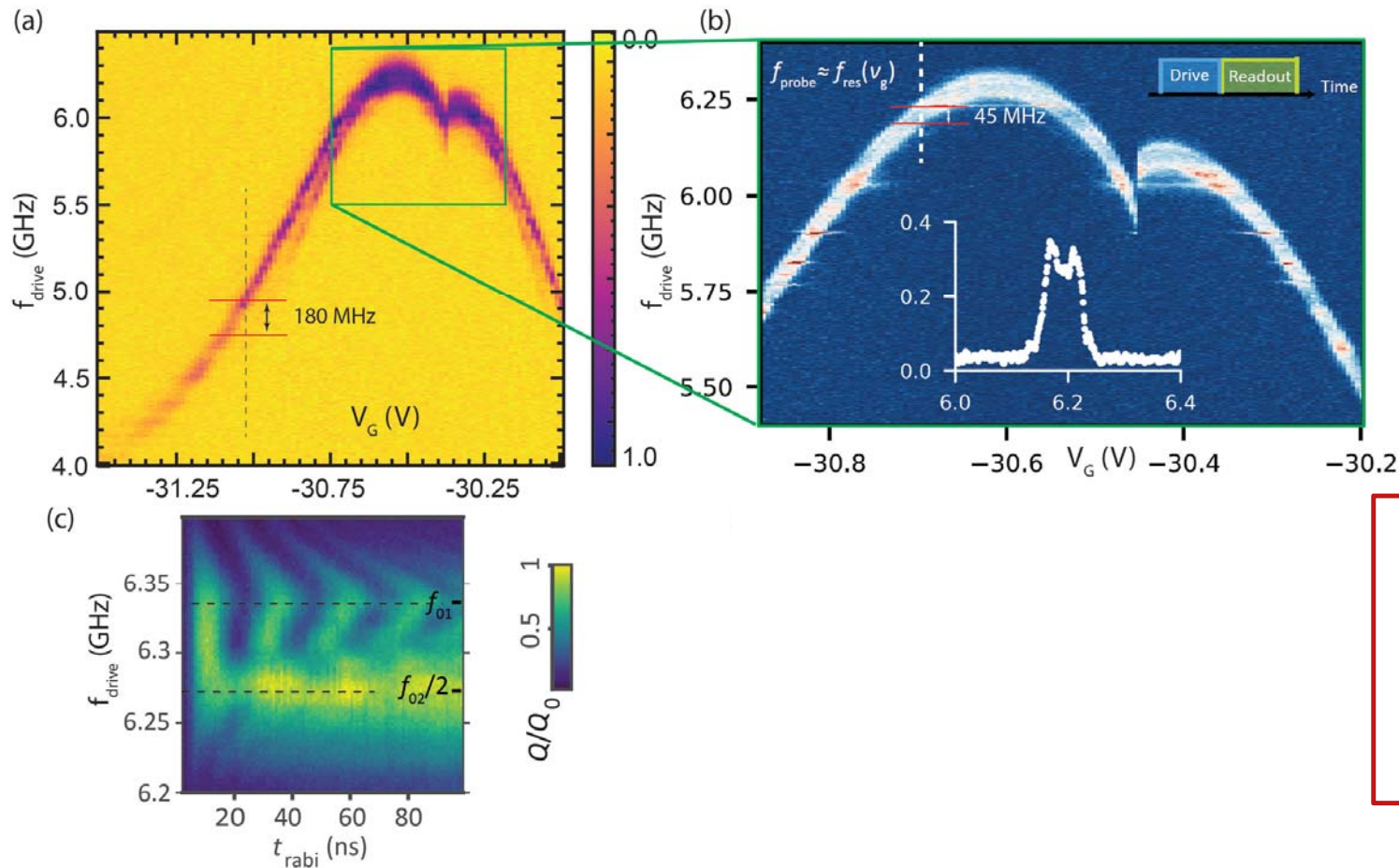
$$g/2\pi \cong 47 \text{ MHz}$$

GeSi core-shell nanowire Gatemon



Extracted the **anharmonicity α** from second peak in **two-tone spectroscopy**; supported by additional signal in Rabi measurements

GeSi core-shell nanowire Gatemon



Extracted the **anharmonicity α** from second peak in **two-tone spectroscopy**; supported by additional signal in Rabi measurements

Using measured values of f_q and α , we can make some predictions, shown in (d):

Josephson junction is in the QPC limit with $N=2$ channel of which one channel has (close to) unit transmission

GeSi core-shell nanowire Gatemon

A. Kringhoj et al., Anharmonicity of a superconducting qubit with a few-mode Josephson junction, Phys. Rev. B 97, 060508 (2018).

Effective Hamiltonian *PRB* 97, 060508(R) (2018)

$$\hat{H} = 4E_C \hat{n}^2 + \sum_i \Delta \sqrt{1 - T_i} \sin^2(\hat{\phi}/2)$$

$$(f_{12} - f_{01})/h = \alpha \approx -E_C \left(1 - \frac{3 \sum_i T_i^2}{4 \sum_i T_i} \right)$$

JJ potential in short-junction limit

simulation $\rightarrow E_C/h \cong 250$ MHz

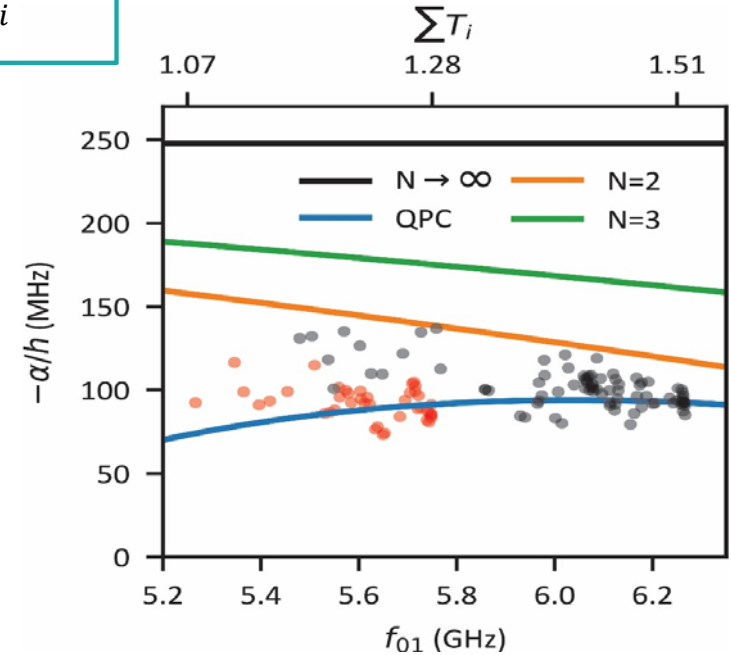
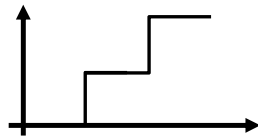
$$hf_{01} \cong \sqrt{8E_C E_J}$$

$$E_J = \frac{\Delta}{4} \sum T_i$$

experiment yields f_{01} (and α)
from E_J one gets $\sum T_i$

N=2 Assuming N channels with equal transmission:
N=3 $E_J = \frac{\Delta}{4} N \cdot \bar{T}$

QPC Assuming steplike channels:

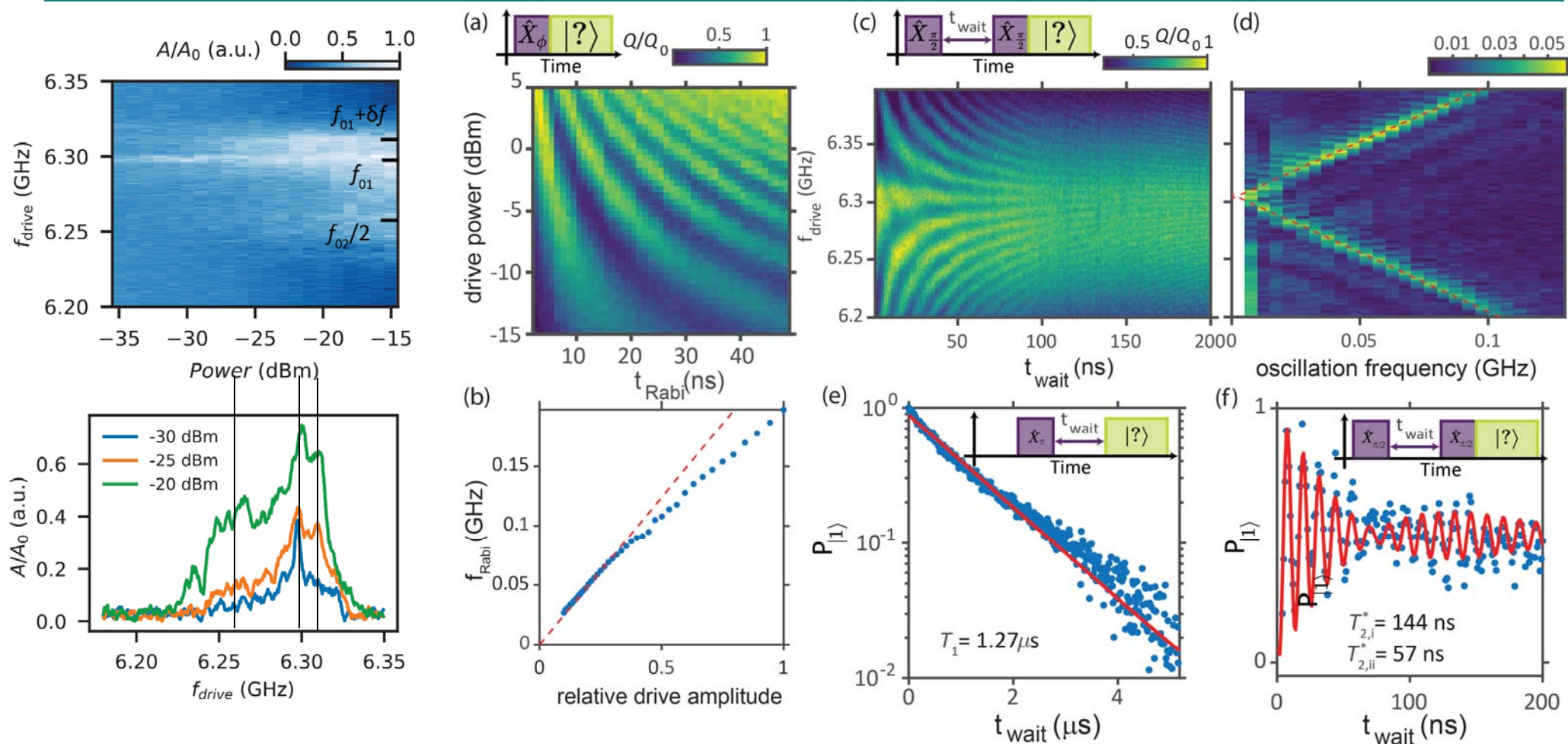


conventional transmon **Low transparent many channels**

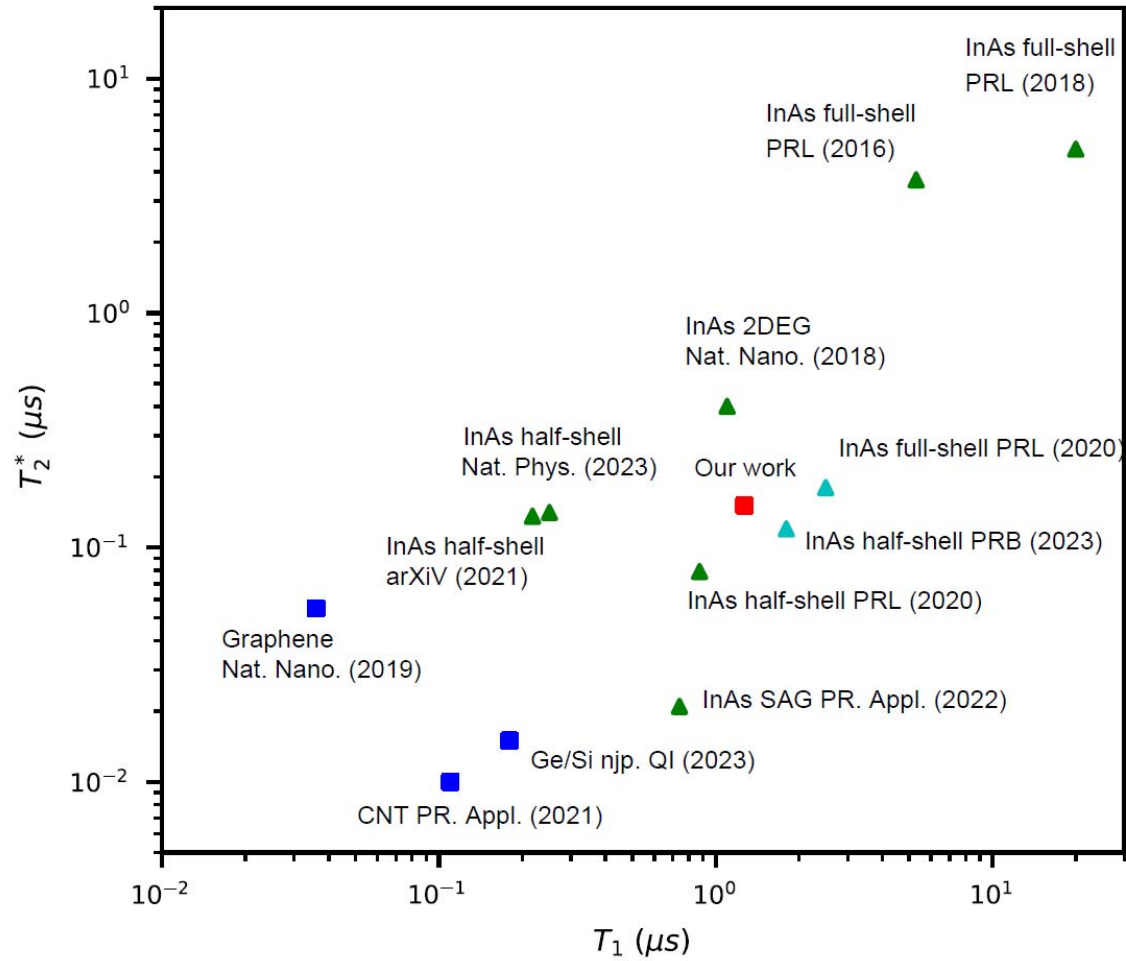


Ge/Si NW gatemon **High transparent 1-2 channels**

GeSi core-shell nanowire Gatemon



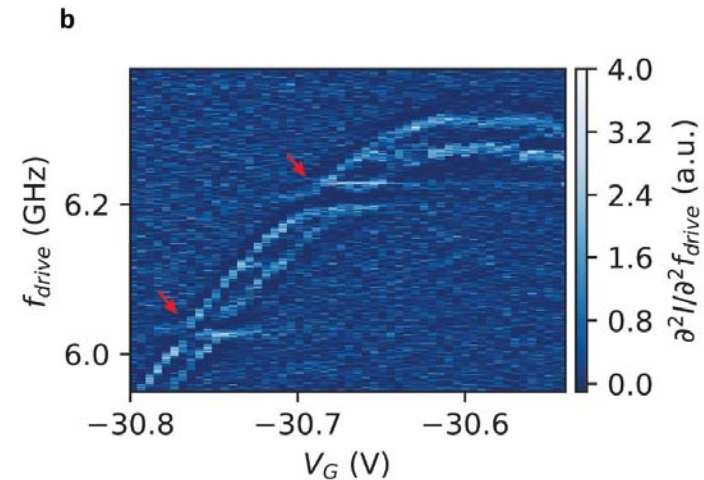
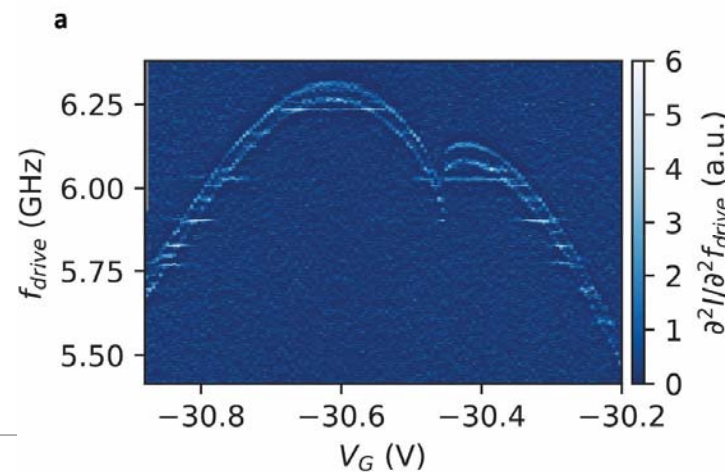
GeSi core-shell nanowire Gatemon



GeSi core-shell nanowire Gatemon

Conclusions:

- gatemon in short strained Ge channel
- at most 2 channels (most likely one that dominates the physics)
- interface from **Al to Ge channel is highly transmissive**
- up to a factor four **reduced anharmonicity**
- relaxation and coherence comparable to state-of-the-art III-V gatemons
- **ideal for Andreev spin qubits**
- two-level fluctuators; needs cleaner processing and etching of Si-oxide



Thank you
for your attention.